

Top Quark Physics

(Lecture II)

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Outline

- Production and decays
- Pair production cross section
- Top quark mass
- Event kinematics

China Center of Advanced Science & Technology
Beijing, China
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Building Blocks of Matter

Quarks	u	c	t
	up	charm	top
Leptons	d	s	b
	ν_e e- Neutrino	ν_μ μ - Neutrino	ν_τ τ - Neutrino
	e electron	μ muon	τ tau
	I	II	III
The Generations of Matter			

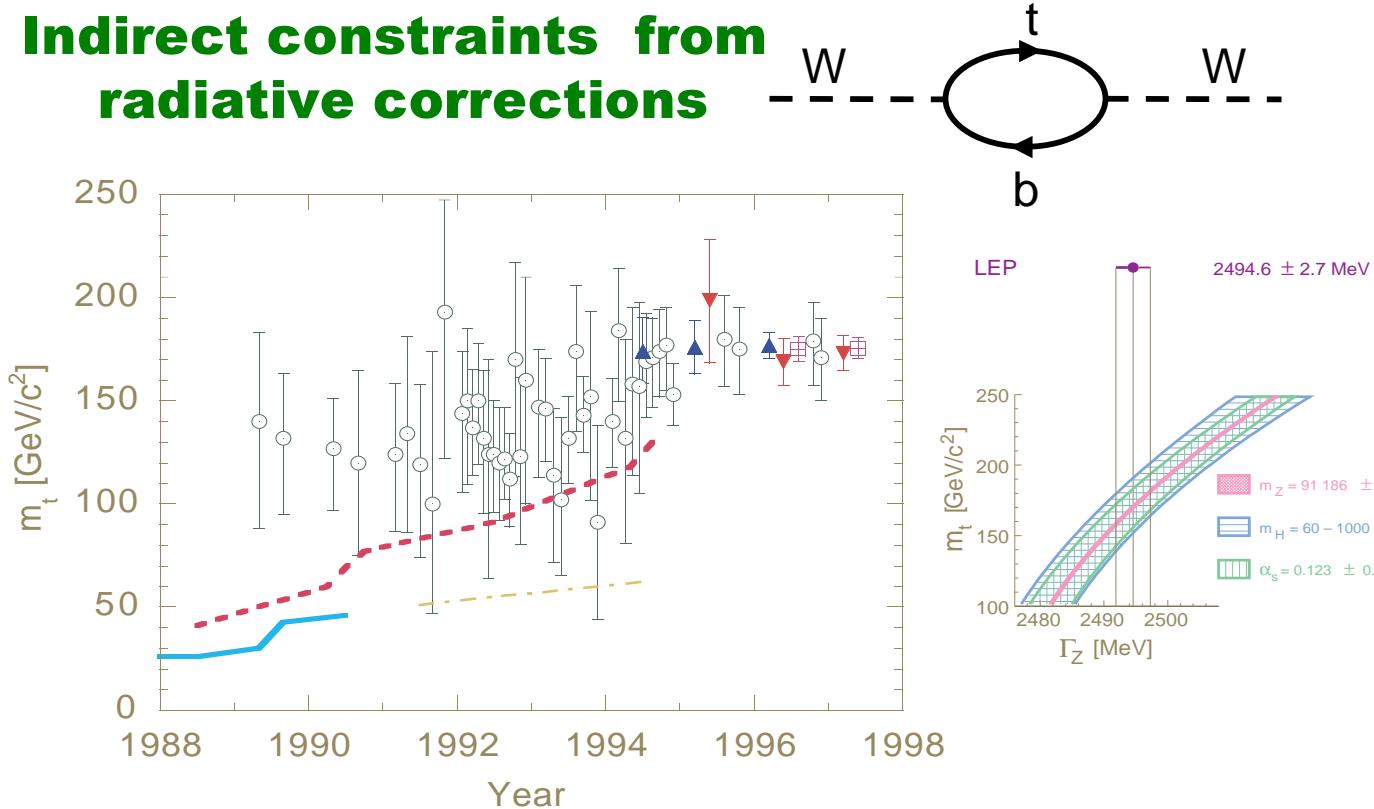
- **1975 – discovery of the τ lepton**
Perl et al, PRL 35, 1489 (1975)
- **1977 – discovery of the b quark**
Herb et al, PRL 39, 252 (1977)
 - weak isospin = -0.504
 - no flavor changing neutral current
 - need weak isospin partner
- **1995 – discovery of the top quark**
CDF, PRL 74, 2626 (1995)
DØ, PRL 74, 2632 (1995)
- **2000 – observation of tau neutrinos**

Hunt for the Top Quark

Direct searches at colliders

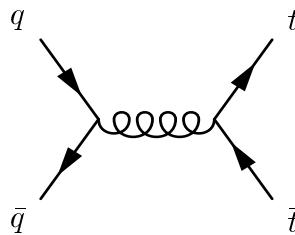
1979-84	PETRA (DESY)	e^+e^-	$m_{top} > 23.3 \text{ GeV}$
1987-90	TRISTAN (KEK)	e^+e^-	$m_{top} > 30.2 \text{ GeV}$
1989-90	SLC (SLAC) LEP (CERN)	e^+e^-	$m_{top} > 45.8 \text{ GeV}$
1990	SppS (CERN)	pp	$m_{top} > 69 \text{ GeV}$
1991	Tevatron (FNAL)	pp	$m_{top} > 77 \text{ GeV}$
1992	Tevatron (FNAL)	pp	$m_{top} > 91 \text{ GeV}$
1994	Tevatron (FNAL)	pp	$m_{top} > 131 \text{ GeV}$

Indirect constraints from radiative corrections

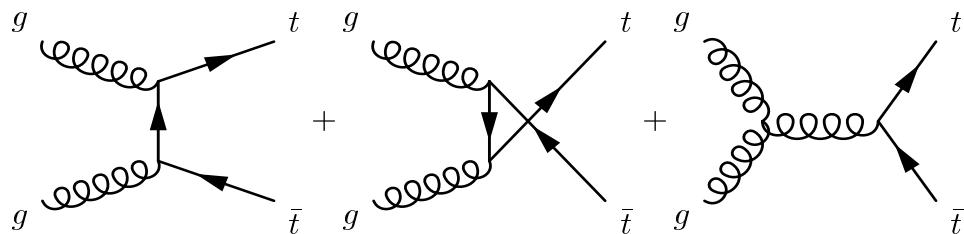


Top Quark Production

Top-antitop quark pair production

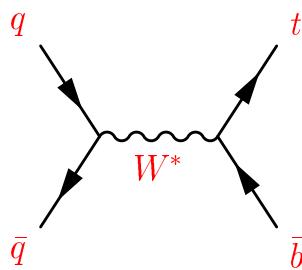


$$q\bar{q} \rightarrow t\bar{t}$$

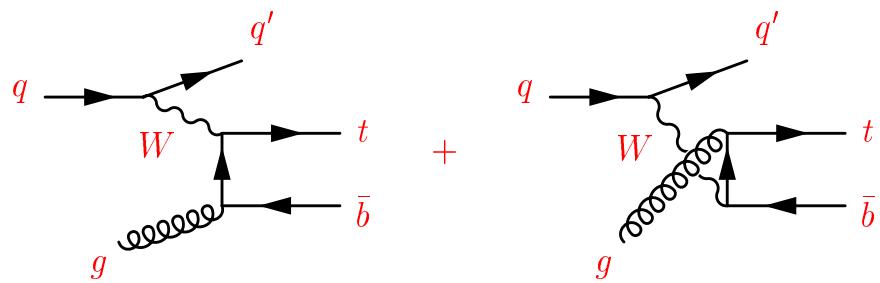


$$gg \rightarrow t\bar{t}$$

Single top quark production



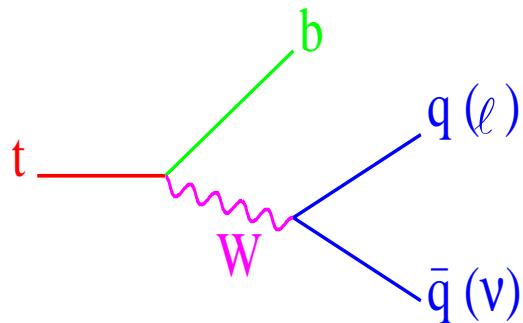
$$q\bar{q} \rightarrow t\bar{b}$$



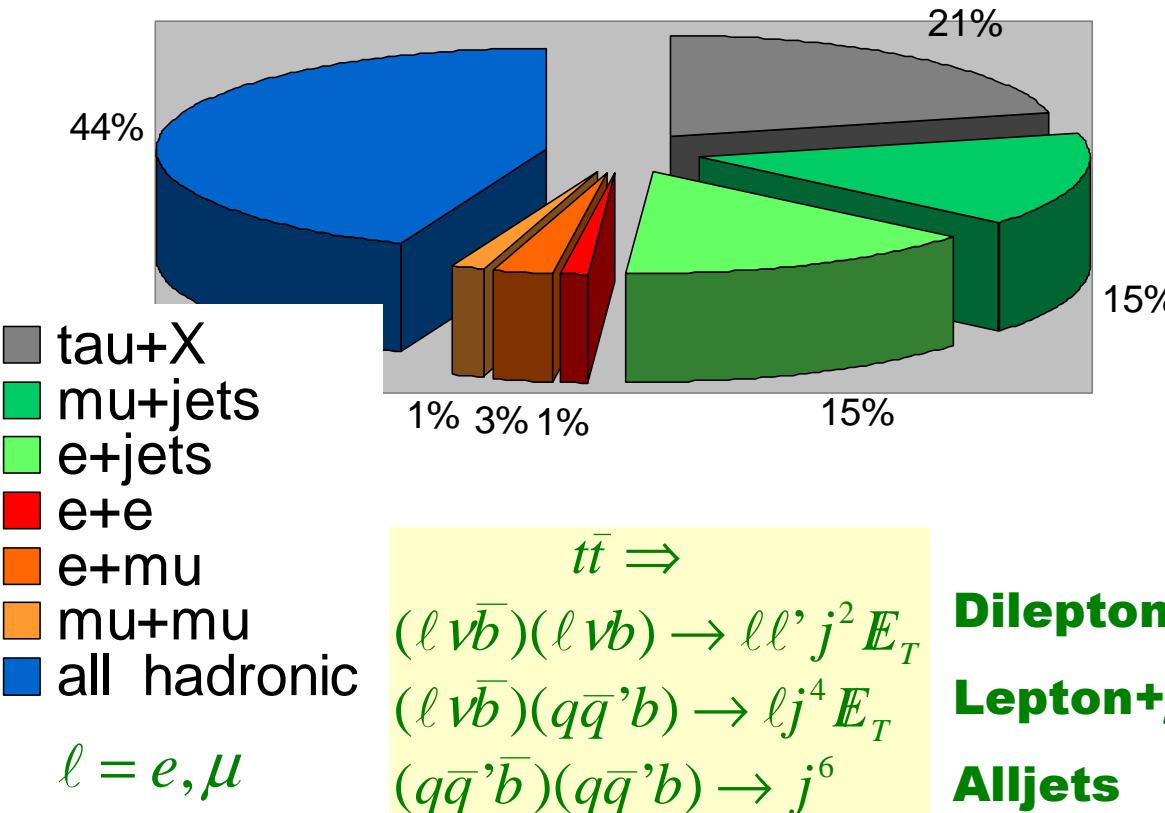
$$qg \rightarrow q' t\bar{b}$$

Top Quark Decay

In the standard model,
the top quark is short lived and decay
almost exclusively to W and b quark

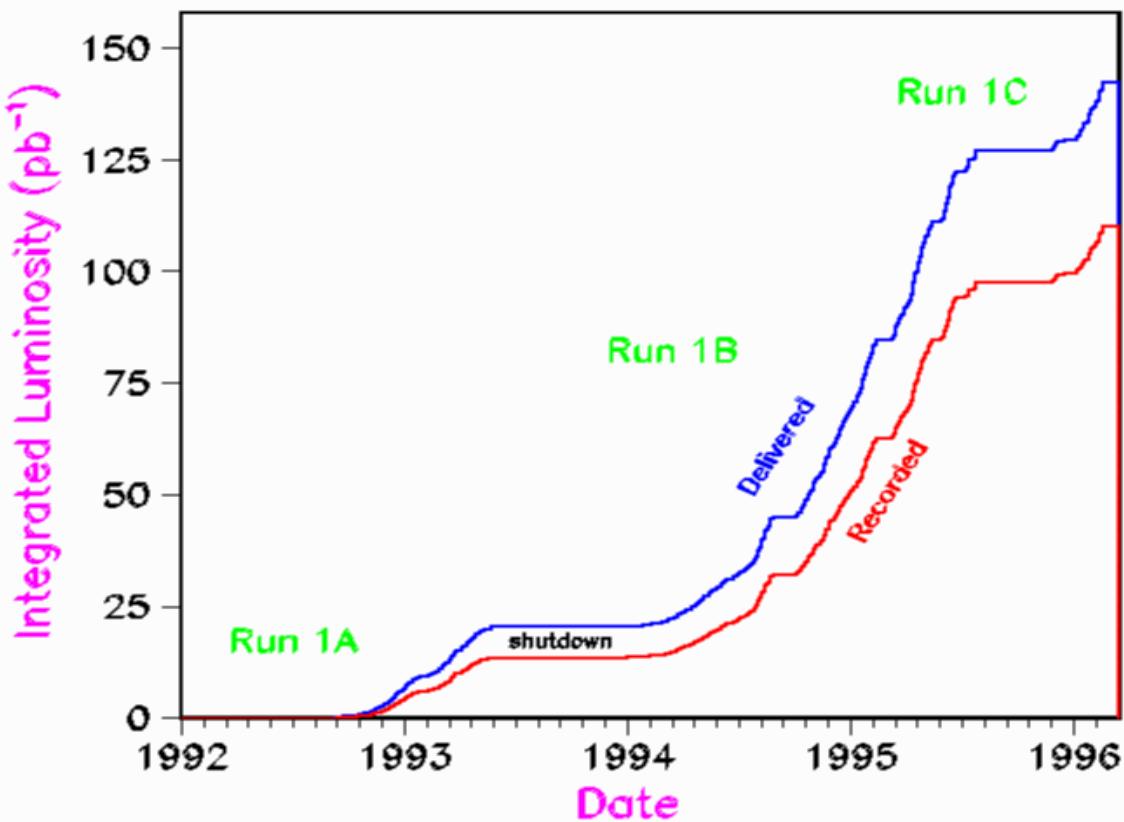


$(W \rightarrow \ell\nu, W \rightarrow \ell\nu) \Rightarrow \ell\ell$
 $(W \rightarrow \ell\nu, W \rightarrow qq') \Rightarrow \ell + \text{jets}$
 $(W \rightarrow qq', W \rightarrow qq') \Rightarrow \text{All-jet}$



Data Sample

- Tevatron Run I (1992–1996):
accumulated $\sim 120 \text{ pb}^{-1}$ integrated luminosity
at a center-of-mass energy of 1.8 TeV



- Tevatron Run II (2001–2006):
expect to accumulate $\sim 20 \text{ fb}^{-1}$ (2 fb^{-1} Run IIa)
at 2.0 TeV

Key for top quark physics:

- good lepton identification
- good missing E_T resolution
- efficient b-jet tagging capability

Dilepton Final State

$$t\bar{t} \rightarrow (\ell\nu\bar{b})(\ell'\nu b) \rightarrow \ell\ell' j^2 E_T$$

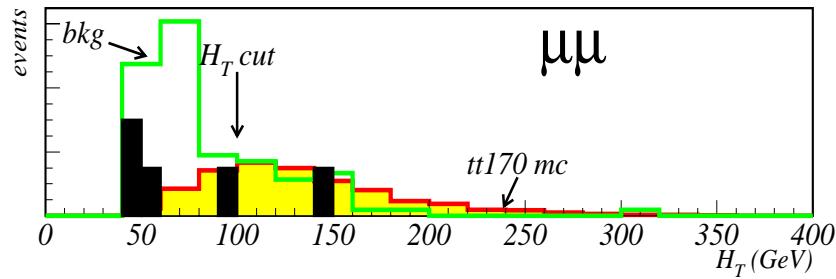
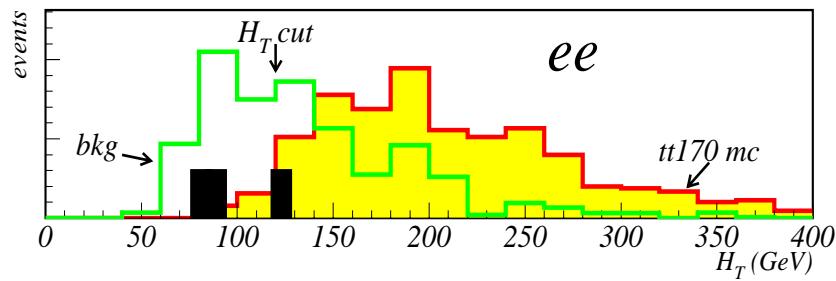
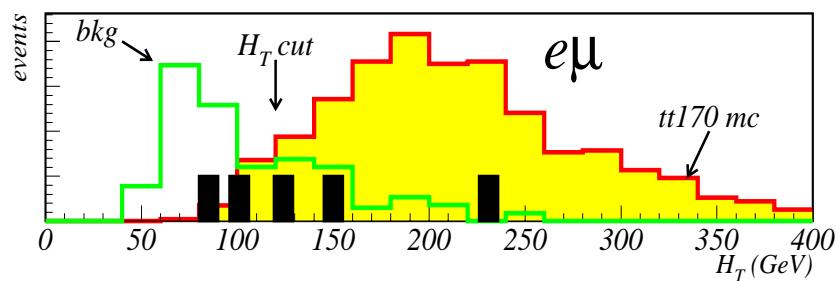
Small rate (5%) but lowest backgrounds

Characteristics:

- two isolated high pT leptons
- 2 jets from b-quarks
- significant missing transverse energy

Major backgrounds:

- $Z(\rightarrow \ell\ell) + \text{jets}$, WW , $Z \rightarrow \tau\tau$
- **QCD**: $b\bar{b}, c\bar{c} \rightarrow \ell\ell$
- **Instrumental: misidentification**



Lepton+Jets Final State

$$t\bar{t} \rightarrow (\ell\nu b)(q\bar{q}'b) \rightarrow \ell j^4 E_T$$

Moderate rate (15%) and reasonable backgrounds

Characteristics:

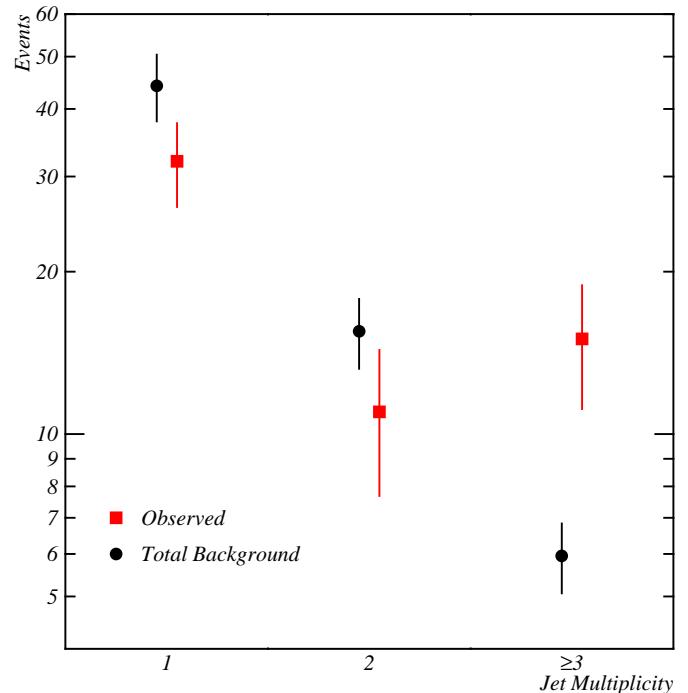
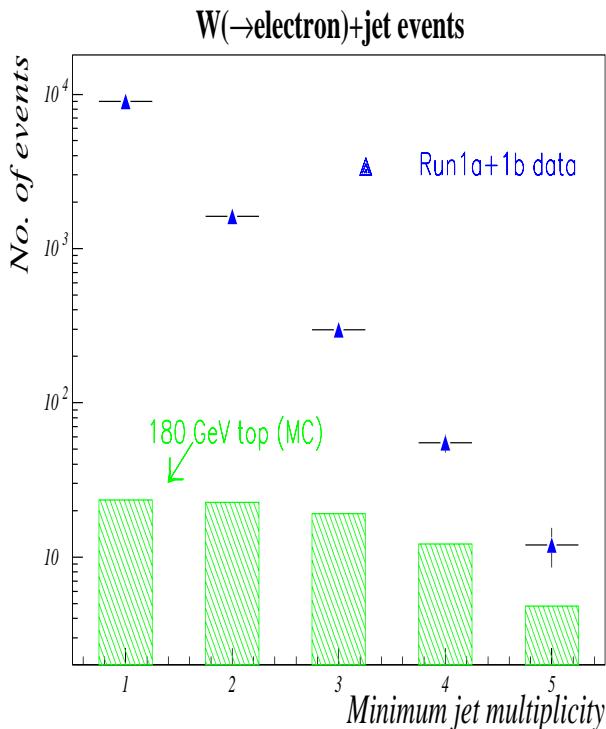
- One isolated high pT lepton
- large transverse momentum imbalance
- multiple jets

Major backgrounds:

- W+jets production
- Instrumental:
misidentification and mismeasurement

Strategies:

- Explore the difference in topologies
- b-jets identification through ($b \rightarrow \mu$) decays



Alljets Final State

$$t\bar{t} \rightarrow (q\bar{q}'\bar{b})(q\bar{q}'b) \rightarrow j^6$$

Largest rate (44%) but highest backgrounds

Characteristics:

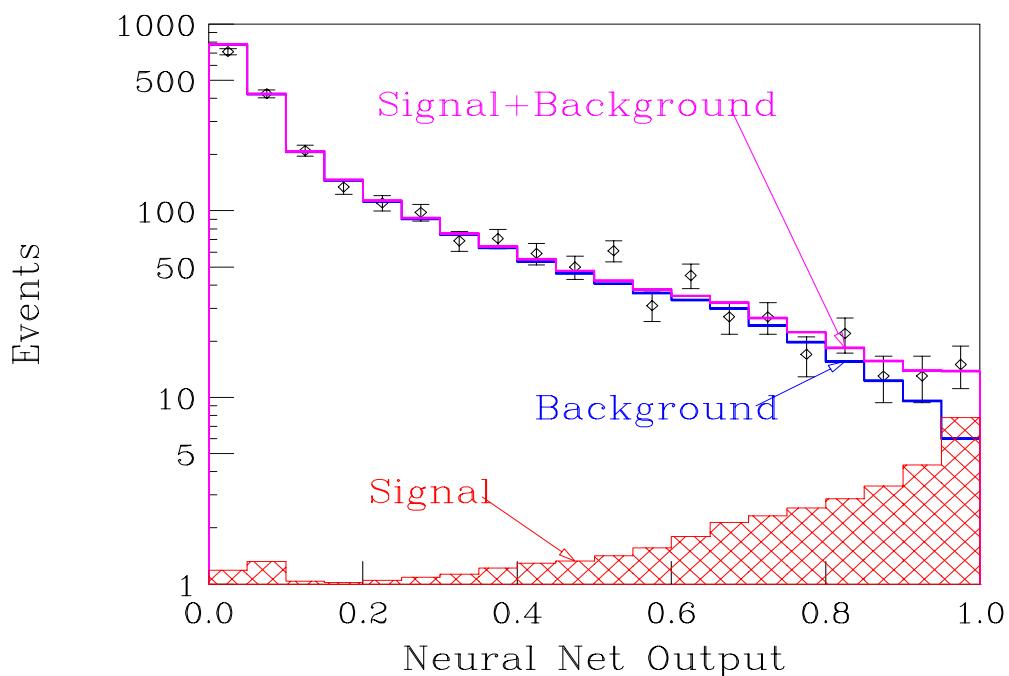
- Six or more jets (two b-jets)
- No high pT lepton nor large transverse momentum imbalance

Backgrounds:

- Huge QCD multi-jet production (**S/B=1/2000 on tape**)

Strategies:

- Tag b-jets using $b \rightarrow \mu$ decays
- Neural network to separate signal from backgrounds



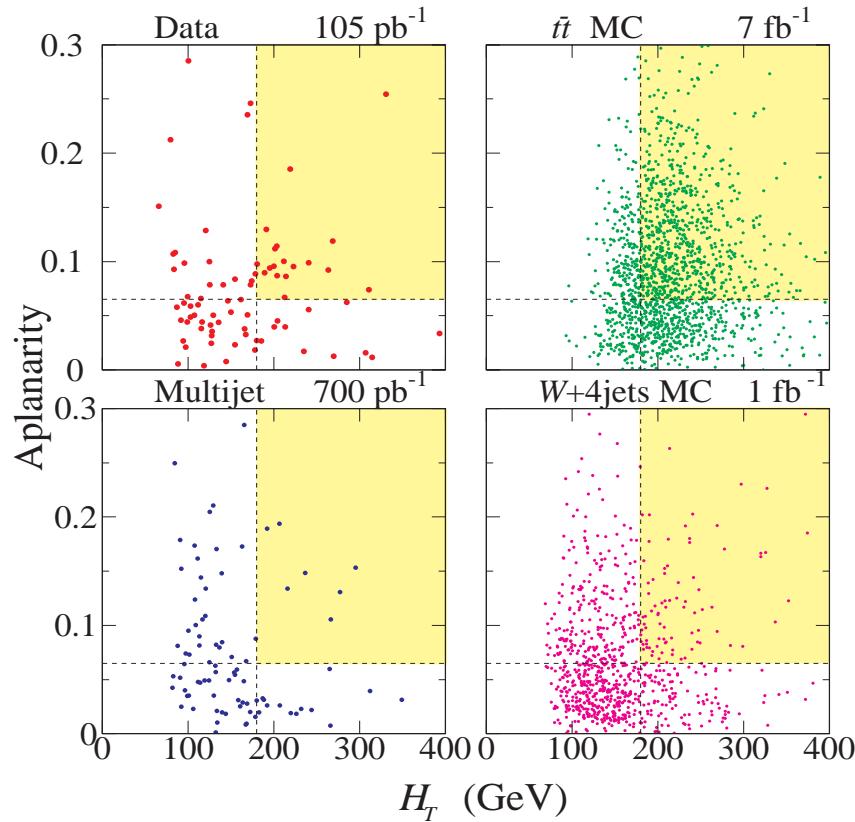
Event Selection

Top quark pair events are topologically different from most of the background events

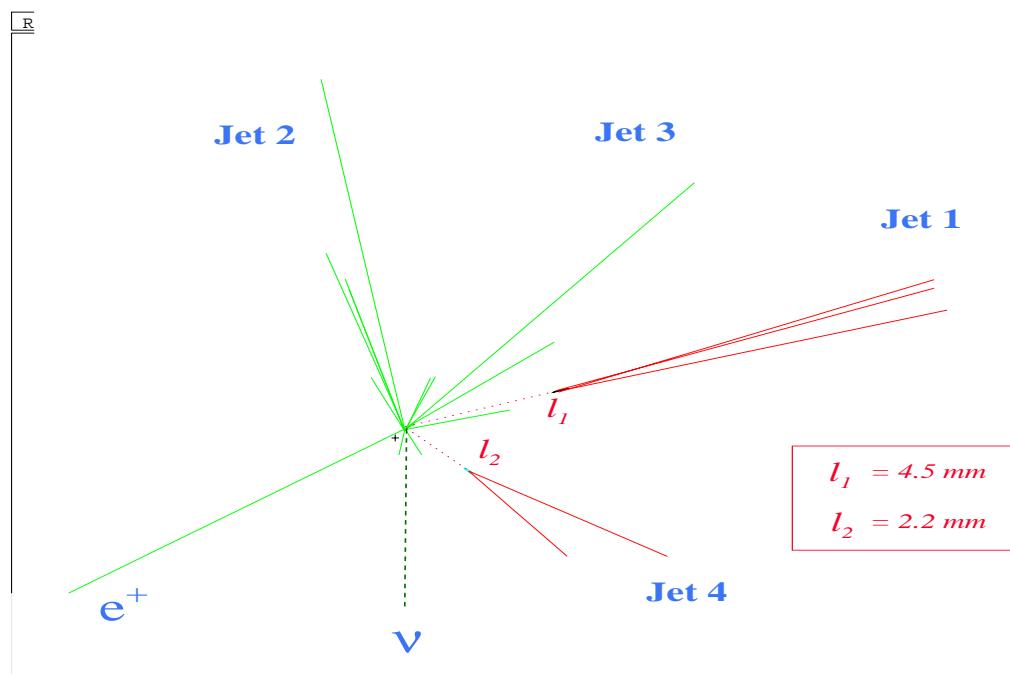
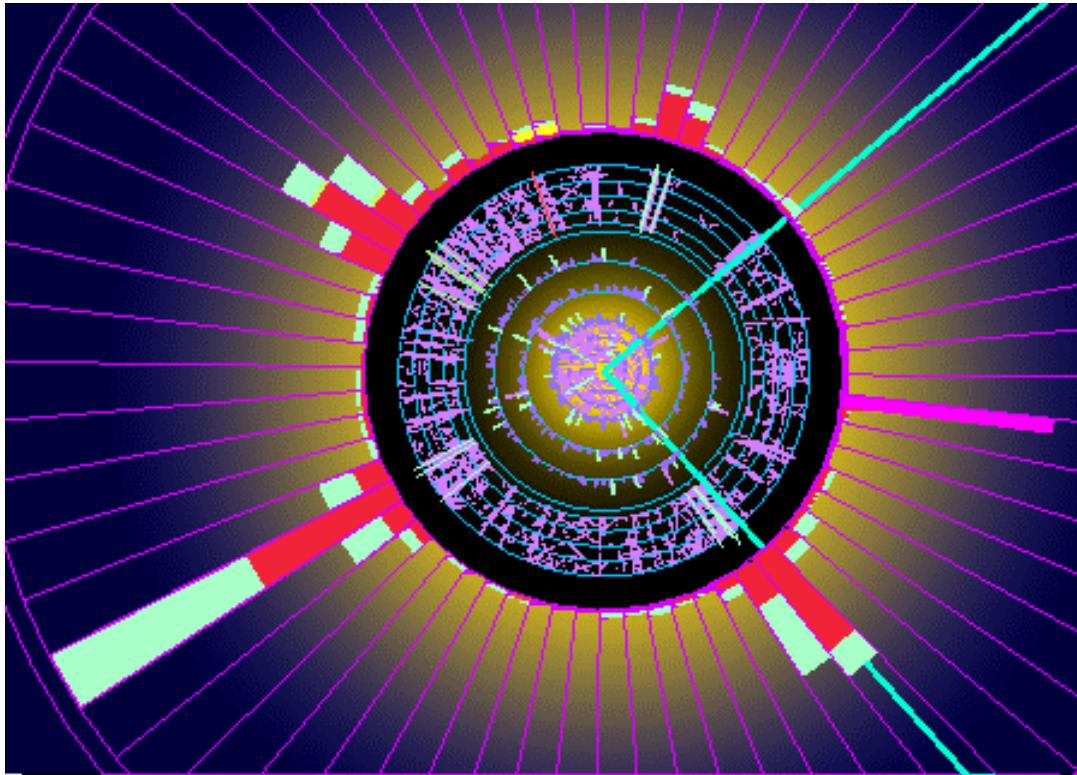
- top quark events are more spherical
- every event has two jets from b-quarks
- leptonic decays of W bosons:
 - high pT isolated leptons
 - large missing transverse momentum

Depending on the final state, sophisticated selection criteria were developed to distinguish signal events from backgrounds

Selection of $\ell + \text{jets}$ events



Candidate Events



$M_{\text{top}}^{\text{Fit}} = 170 \pm 10 \text{ GeV}/c^2$

24 September, 1992
run #40758, event #44414

Selection Criteria

Dilepton Final State

		$e\mu$	ee	$\mu\mu$
Lepton	E_T (GeV)	≥ 15	≥ 20	≥ 15
	$ \eta $	<2.5 (e) <1.7 (μ)	<2.5	<1.7
Missing E_T (GeV)		≥ 20	≥ 25	-
Jets	#	≥ 2	≥ 2	≥ 2
	E_T (GeV)	≥ 20	≥ 20	≥ 20
	$ \eta $	<2.5	<2.5	<2.5
H_T (GeV)		≥ 120 (jets+e)	≥ 120 (jets+e ₁)	≥ 100 (jets)

Lepton+Jets Topological Analysis

Lepton	E_T (GeV)	>20	>20
	$ \eta $	<2	<1.7
Missing E_T (GeV)		>25	>20
2 Jets	E_T (GeV)		>15
	$ \eta $		<2
Shape	$E_T^\ell + E_T$		>60 GeV
	$\sum E_T^{jets}$		>180 GeV
	Aplanarity		>0.065
Others		$ \eta_W < 2$	
		Veto μ -tagged events	

Selection Criteria

Lepton+Jets Tagged Analysis

- **two b-quark jets in each signal events**
- **background events have minimal heavy flavor content**
- **tag b-jet by presence of μ in jet**
 - **~20% of the signal events have a detectable μ associated with a jet**
 - **~2% of W+jets background events have a μ associated with a jet**
- **relax shape cuts to require only 3 jets**

Event Selection

		e+jets/ μ	μ +jets/ μ
Lepton	E_T (GeV)	>20	
	$ \eta $	<2	<1.7
Missing E_T (GeV)		>20	
$Z \rightarrow \mu\mu$ fit		–	$\text{Prob}(\chi^2) < 0.01$
Soft μ		$p_T > 4 \text{ GeV}$ and $\Delta R(\mu, \text{jet}) < 0.5$	
Jets		≥ 3 jets, $E_T > 20 \text{ GeV}$, $ \eta < 2$	
Shape		$\text{Aplanarity} > 0.04$	
		$H_T > 110 \text{ GeV}$	

Event Sample

$$t\bar{t} \rightarrow (\ell\nu b)(\ell'\nu b) \rightarrow \ell\ell' j^2 E_T$$

	DØ	CDF	
	ee, e _{μ, μμ} , (e _ν)	e _{μ, μμ}	e _{τ, μτ}
data	5 (4)	9	4
background	1.4±0.4 (1.2±0.4)	2.4±0.5	2.0±0.4

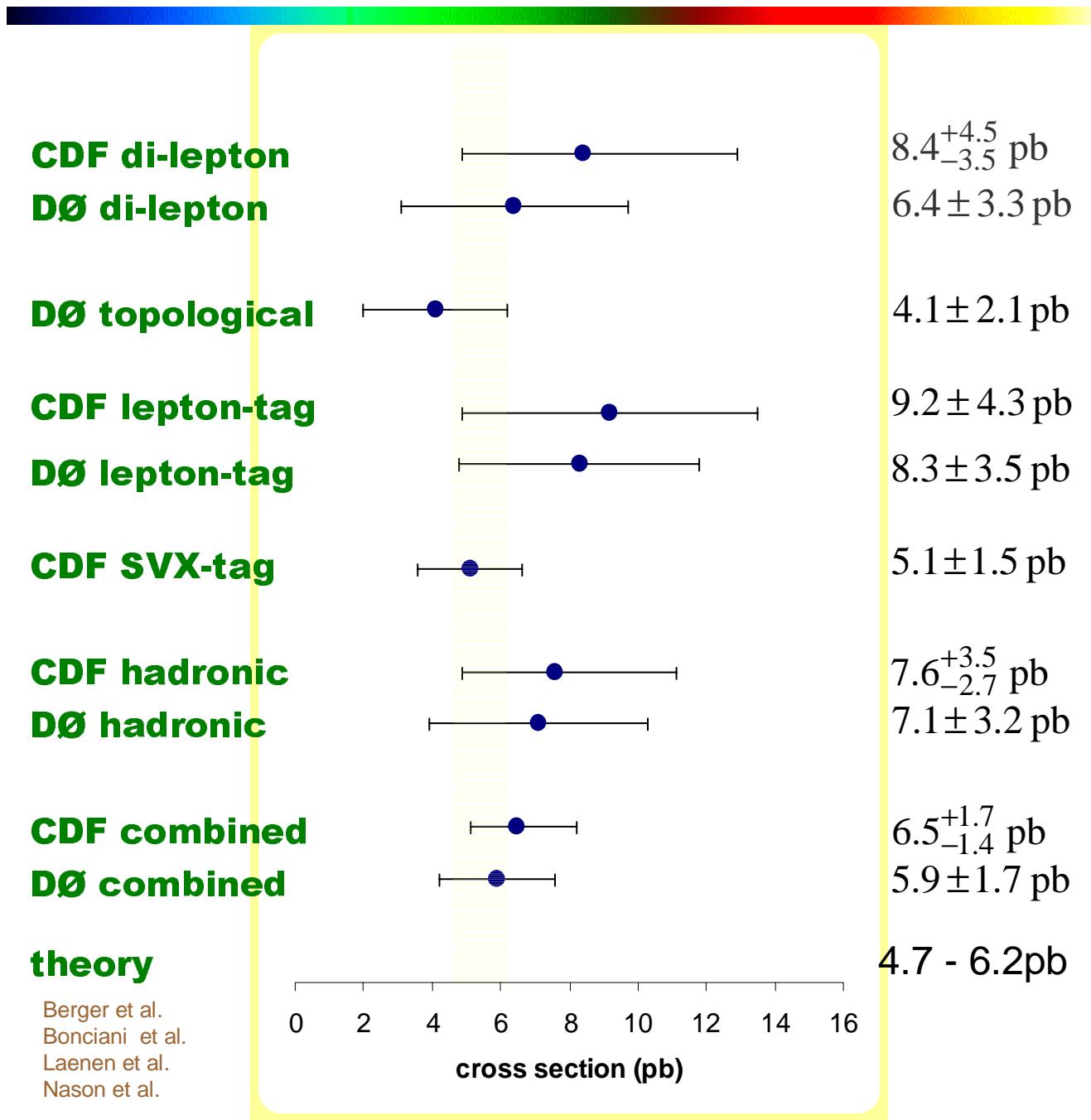
$$t\bar{t} \rightarrow (\ell\nu b)(q\bar{q}'b) \rightarrow \ell j^4 E_T$$

	DØ		CDF	
	topological	lepton-tag	SVX-tag	lepton-tag
data	19	11	29	25
background	8.7±1.7	2.4±0.5	8.0±1.0	13.2±1.2
		11 events in common		

$$t\bar{t} \rightarrow (q\bar{q}'\bar{b})(q\bar{q}'b) \rightarrow j^6$$

	DØ	CDF
data	41	187
background	24±2.4	151±10

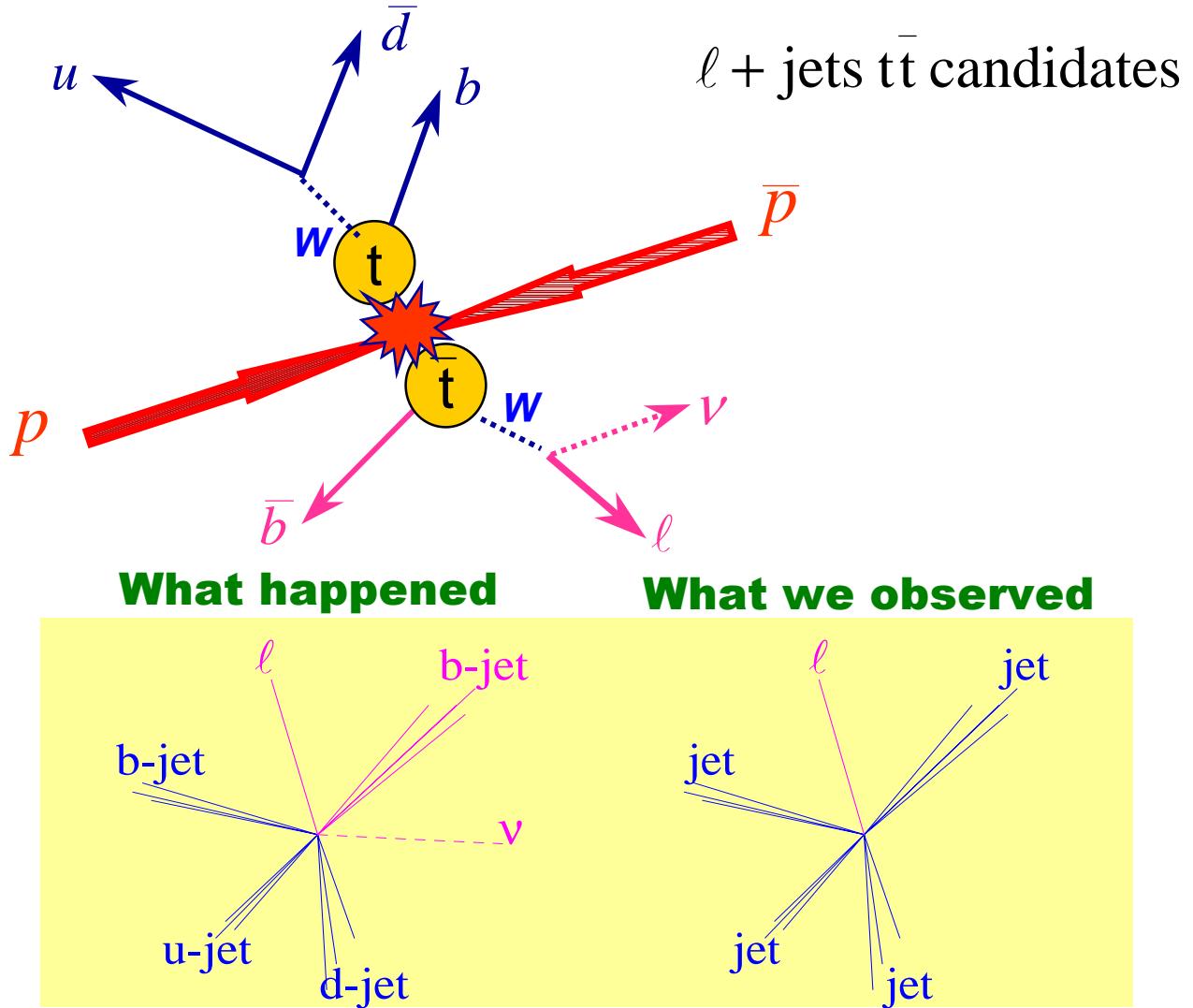
Cross Section



Theoretical uncertainty is about 10% - 20%

Run II precision: ~8% (stat: 4%, syst: 4%, lumi: 4%)

Lepton+Jets Mass Method



For a signal event, we have one unknown (p_z^ν) and three constraints

$$m(\ell\nu) = m(jj) = M_W$$

$$m(j\ell\nu) = m(jjj)$$

Top quark mass can be determined through a 2-constraint fit to event kinematics

Lepton+Jets Mass Method

$$t\bar{t} \rightarrow (\ell\nu b)(q\bar{q}'b) \rightarrow \ell j^4 E_T$$

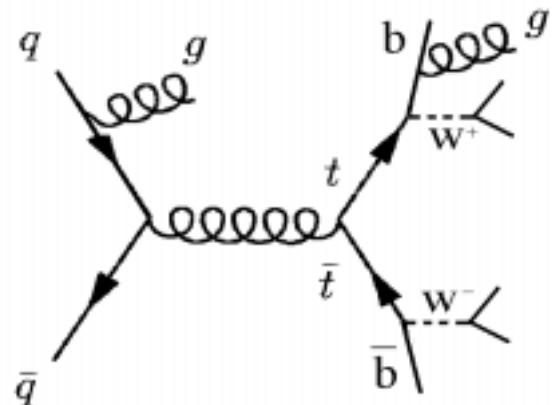
Jet assignment ambiguity

- 12 if no jets is b-tagged
- 6 if one of the jets is b-tagged
- 2 if two jets are b-tagged

Additional complications from

- background events
- detector effect (mismeasurement + resolution)
- initial and final state radiations

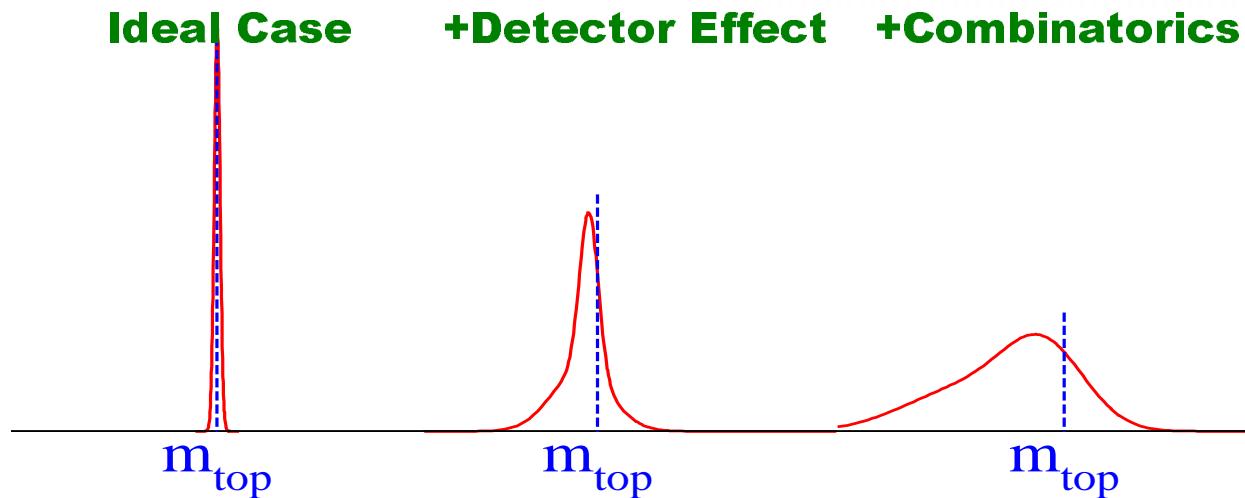
Compare to Monte Carlo
to measure the top
quark mass



Ideal Case

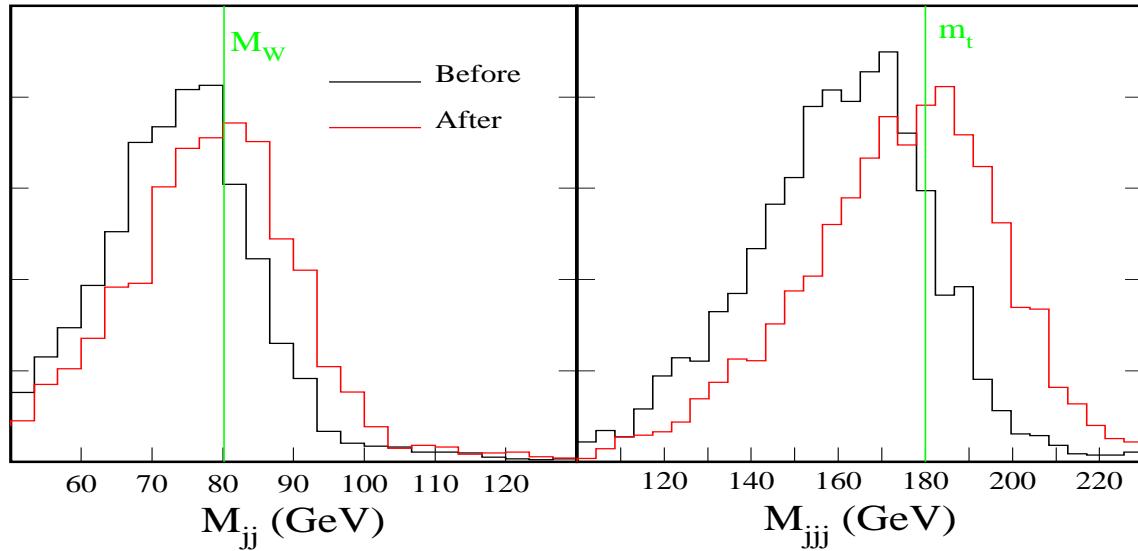
+Detector Effect

+Combinatorics

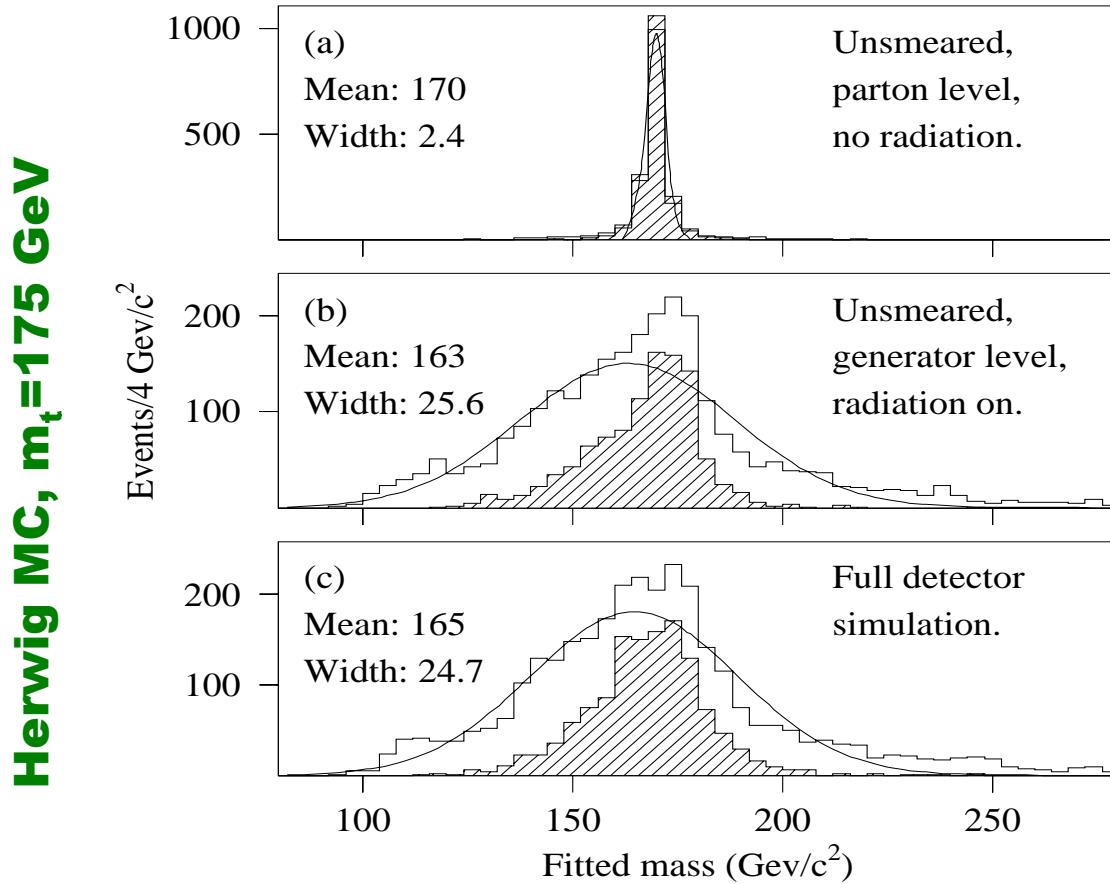


Lepton+Jets Mass Method

Jet Energy Scale



Resolutions and Combinatorics



Lepton+Jets Mass Method

The measured distribution $f(x)$ is compared to those from signal $g_s(x, m_t)$ and from background $g_b(x)$ to extract the top quark mass using a likelihood method

$$L(n_s, n_b, m_t) = \frac{e^{-(n_s+n_b)}}{N!} (n_s + n_b)^N \times \frac{1}{\sqrt{2\pi}\sigma_b} e^{-\frac{(n_b - \langle n_b \rangle)^2}{2\sigma_b^2}} \times \prod_{i=1}^N \frac{n_s g_s(x_i, m_t) + n_b g_b(x_i)}{n_s + n_b}$$



Poisson statistics



Background constraint

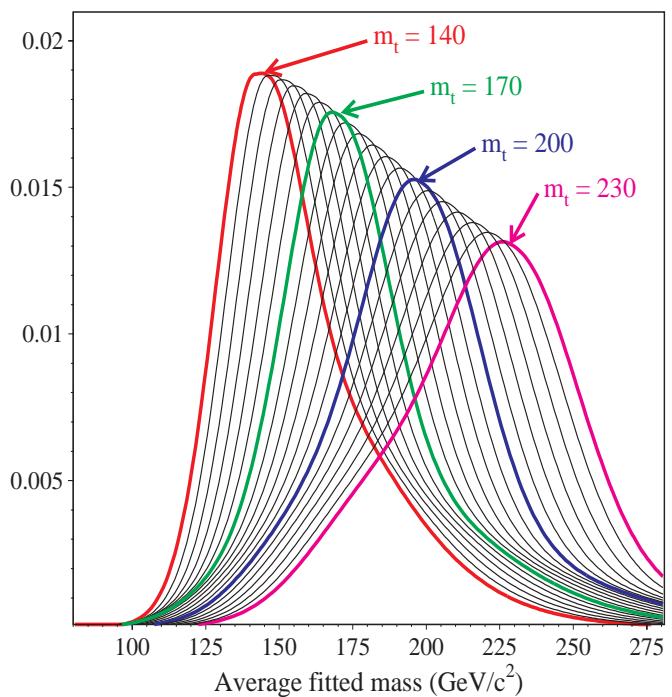
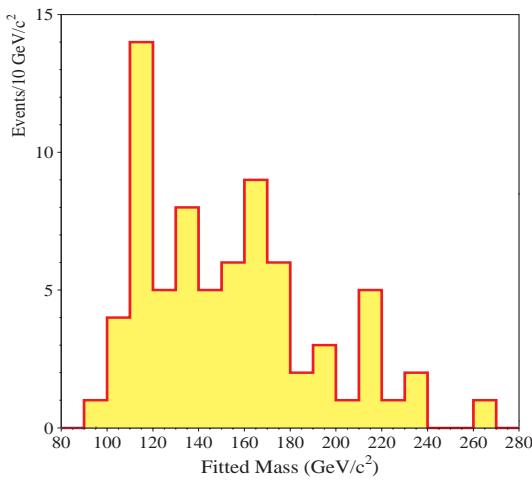


Probability density distribution for the measured variable

n_s = # of signal events

n_b = # of background events

m_t = top quark mass



Multivariate Discriminants

- Some of the cross section selection are mass biased and need to be replaced
- DØ used four-variable multivariate discriminants

$x_1 = E_T$ (transverse momentum imbalance)

$x_2 = A$ (Aplanarity)

$x_3 = \frac{\sum_{i=2}^N E_T^j}{\sum_{i=1}^N E_L^j}$ (centrality)

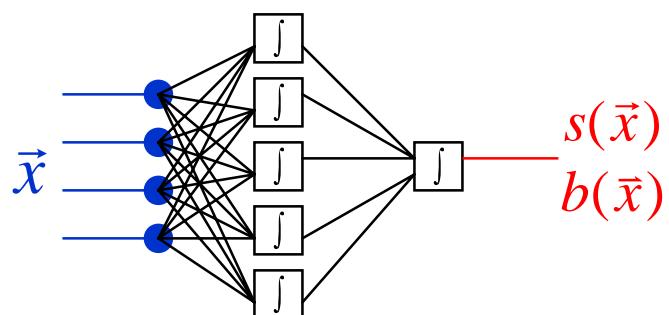
$x_4 = \frac{(\Delta R_{jj}^{\min}) E_T^{\min}}{E_T^L}$ (clusterness)

Likelihood: $L_i(\vec{x}) = \frac{s_i(\vec{x})}{b_i(\vec{x})}$ for each event

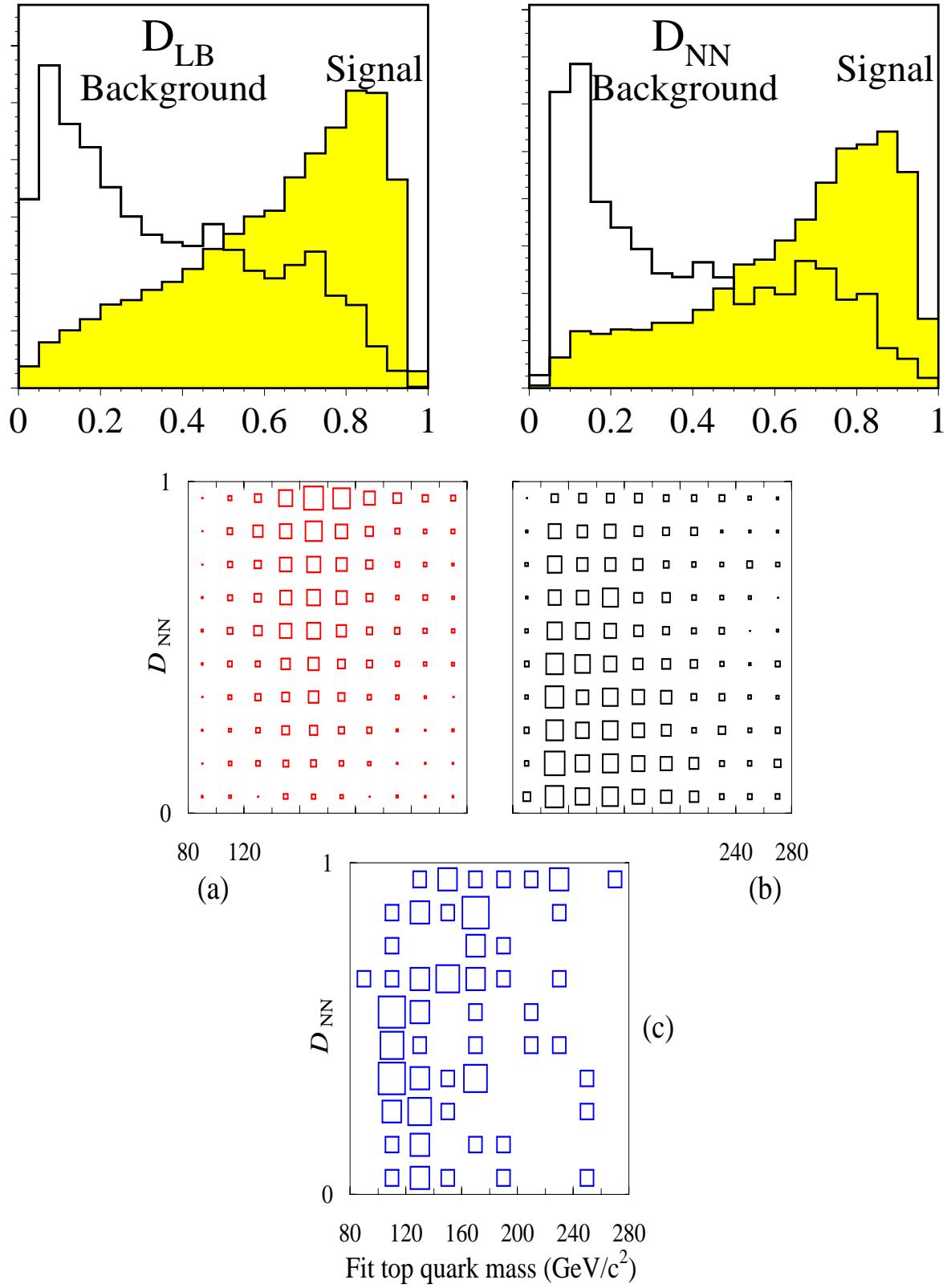
$$D_{LB} = \frac{\prod_{i=1}^N L_i(\vec{x})}{1 + \prod_{i=1}^N L_i(\vec{x})} \Rightarrow \{0 - 1\}$$

Neural Network:

$$D_{NN} = \frac{s(\vec{x})}{s(\vec{x}) + b(\vec{x})} \Rightarrow \{0 - 1\}$$



Multivariate Discriminants



Lepton+Jets Mass Analysis

Analysis Details

$$E_T^\ell > 20 \text{ GeV} \quad |\eta_e| < 2.0 \quad |\eta_\mu| < 1.7$$

$$E_T > 20 \text{ GeV}$$

$E_T^{cal} > 25 \text{ GeV}$ for e + jets; $E_T^{cal} > 20 \text{ GeV}$ for $\mu + \text{jets}$
 ≥ 4 jets, $E_T > 15 \text{ GeV}$, $|\eta_{\text{jet}}| < 2.0$

For events with $b \rightarrow \mu$ tagging

$$p_T^\mu > 4 \text{ GeV} \text{ and } \Delta R = \sqrt{(\delta\eta)^2 + (\delta\phi)^2} < 0.5 \text{ of a jet}$$

For events without $b \rightarrow \mu$ tagging

$$E_T^L \equiv (E_T^\ell + E_T) > 60 \text{ GeV} \text{ and } |\eta_W| < 2.0$$

91 events selected (7 events have $b \rightarrow \mu$ tagged)

Perform kinematic fit on this sample

require $\chi^2 < 10$ and select fit with smallest χ^2

$\Rightarrow 77$ events fitted successfully (5 $b \rightarrow \mu$ tagged events)

Mass from Lepton+Jets

$$t\bar{t} \rightarrow (\ell\nu b)(q\bar{q}'b)$$

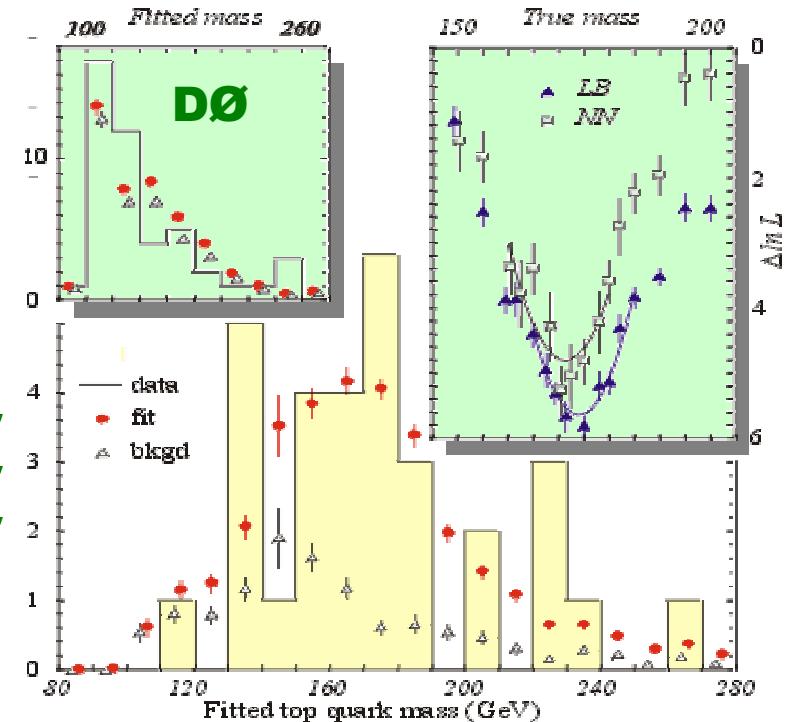
173.3 \pm 5.6 \pm 5.5 GeV

Largest systematics

Jet energy	4.0 GeV
Monte Carlo	3.1 GeV
Noise/pile-up	1.3 GeV

PRL 79, 1197 (1997)

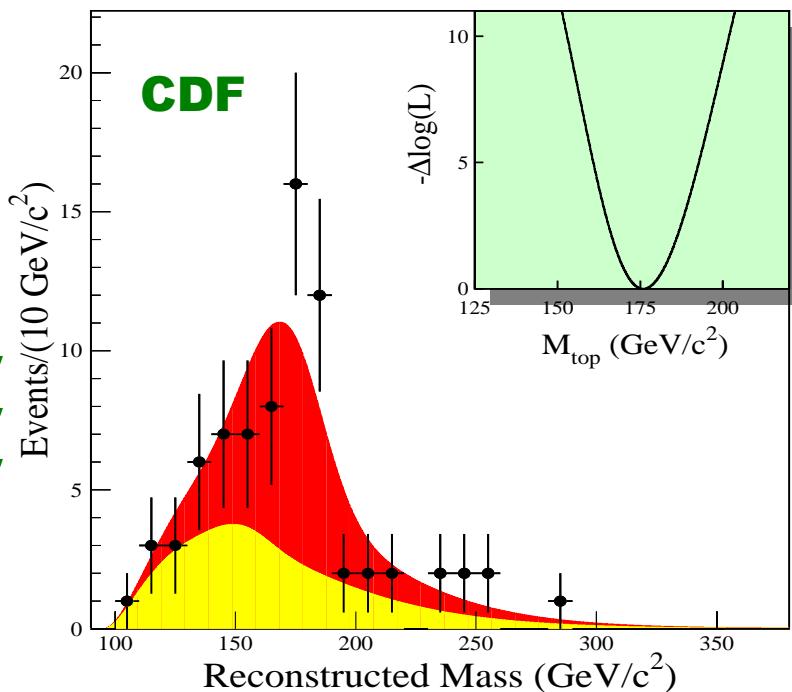
PRD 58, 052001 (1998)



176.1 \pm 5.1 \pm 5.3 GeV

Largest systematics

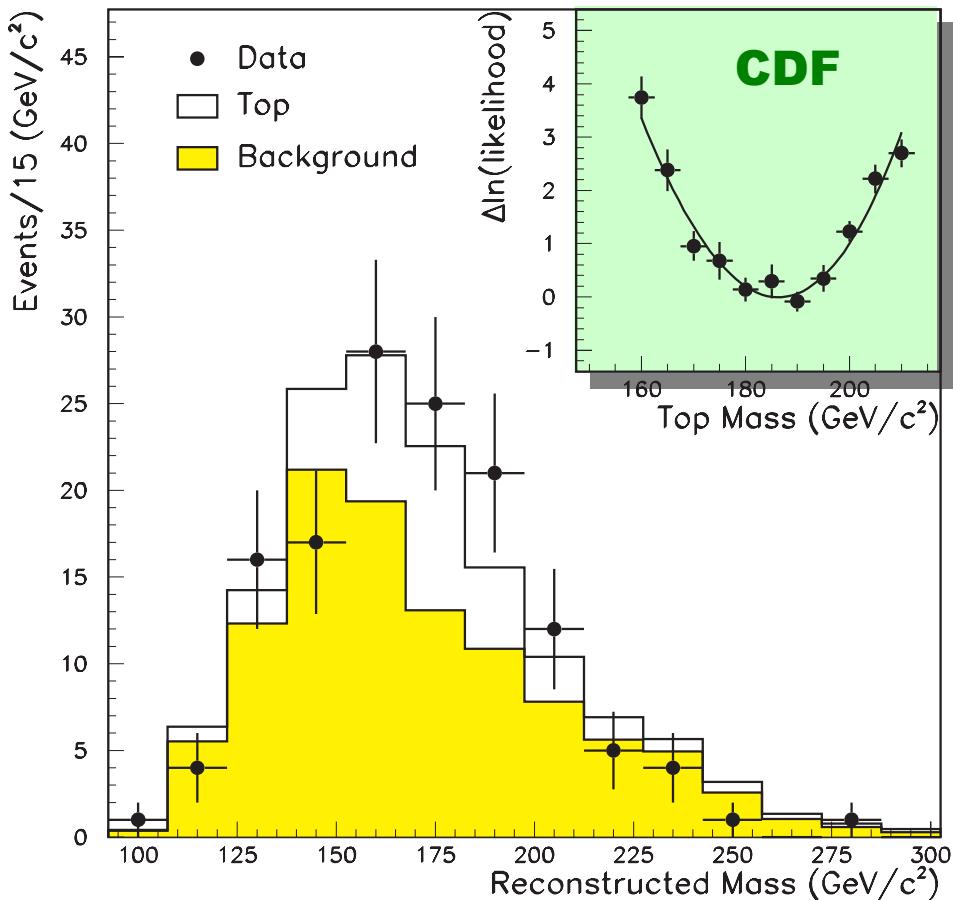
Jet energy	4.4 GeV
Monte Carlo	2.6 GeV
Background	1.3 GeV



Mass from Alljets

$$t\bar{t} \rightarrow (q\bar{q}'b)(q\bar{q}'\bar{b}) \rightarrow j^6$$

m_t can also be determined from alljets final state from a 3-constraint fit



$$m_t = 186.0 \pm 10.0 \pm 5.7 \text{ GeV}$$

Largest systematics

Jet scale	4.4 GeV
Monte Carlo	1.8 GeV
Background	1.3 GeV

Dilepton Mass Method

$$t\bar{t} \rightarrow (\ell\nu b)(\ell'\nu b) \rightarrow \ell\ell' j^2 E_T$$

The final state is underconstrained

- **among the 6 neutrino momentum components, only Σp_T^ν are measured \Rightarrow four unknowns**
- **three constraints:** $m(\ell\nu) = m(\ell'\nu) = M_W$
 $m(\ell\nu b) = m(\ell'\nu b)$

Nevertheless, the final state momentum and angular information is sensitive to the top quark mass.

- **assuming a m_t , calculate a weight for every event characterizing how likely that the event is consistent with the assumed m_t**
- **sum the weight over candidate events**
- **repeat the process for different values of m_t , \Rightarrow weight function $W(m_t)$**
- **the top quark mass is extracted by fitting the observed $W(m_t)$ distribution to those from MC**

K. Kondo, J. Phys. Soc. Jpn. 57, 4126 (1988) and 60, 836 (1991)
R.H. Dalitz and G.R. Goldstein, PRD 51, 4763 (1995).

Dilepton Mass Method

$$t\bar{t} \rightarrow (\ell\nu b)(\ell'\nu b) \rightarrow \ell\ell' j^2 E_T$$

Matrix element and neutrino weightings

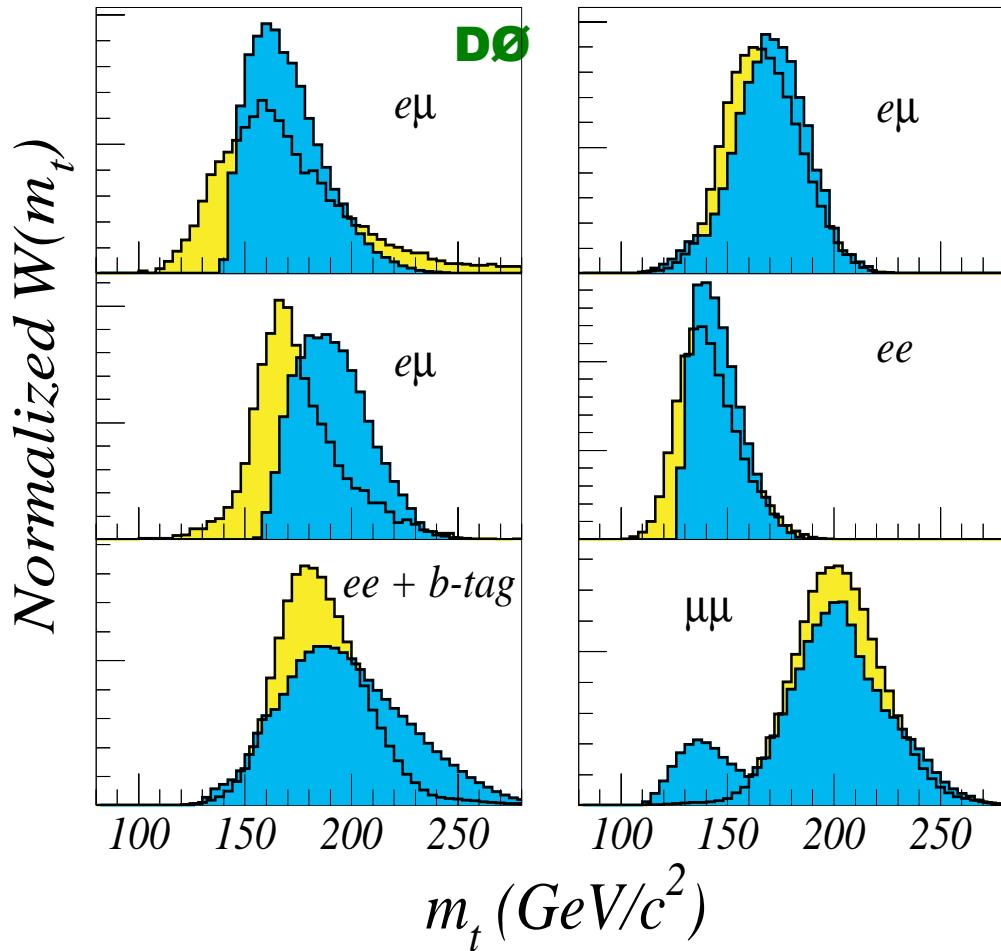
$$w_i^M(m_t) \sim f(x)f(\bar{x})p(E_i^{\ell^+}|m_t)p(E_i^{\ell^-}|m_t)$$

$$w_i^\nu(m_t) \sim \int d\eta d\bar{\eta} g(\eta|m_t)g(\bar{\eta}|m_t)\exp\left[-\frac{(\vec{p}_T - \vec{p}_T^\nu - \vec{p}_T^{\bar{\nu}})^2}{2\sigma^2}\right]$$

$f(x)$: parton distribution function

$p(E^\ell|m_t)$: lepton energy probability for a given m_t

$g(\eta|m_t)$: lepton η probability for a given m_t



Mass from Dilepton

$$t\bar{t} \rightarrow (\ell \nu b)(\ell' \nu b)$$

168.4±12.3±3.6 GeV

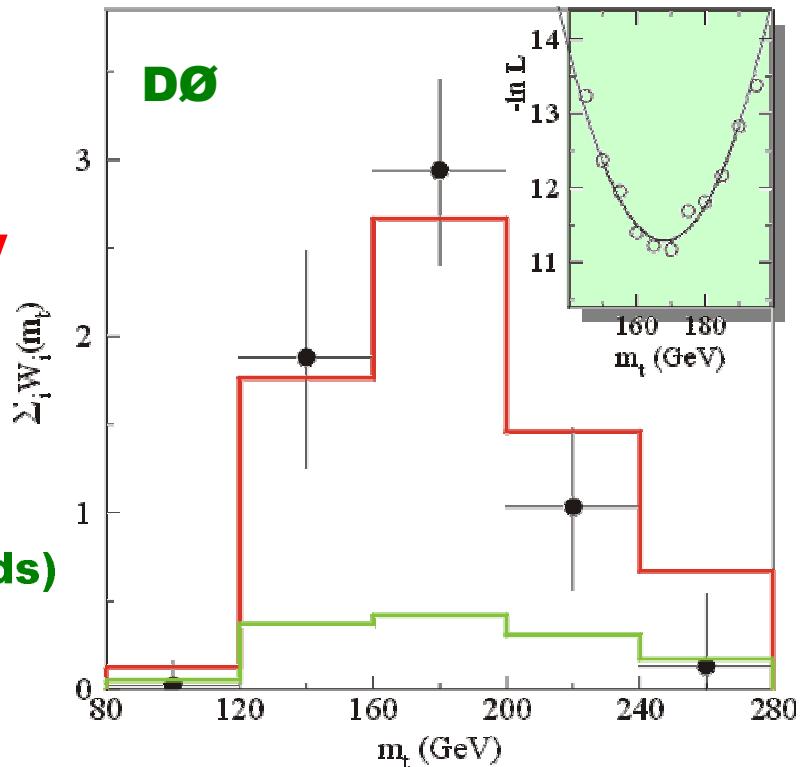
Largest systematics

Jet energy	2.4 GeV
Monte Carlo	1.8 GeV
Noise/pile-up	1.3 GeV

(combining both methods)

PRL 80, 2063 (1998)

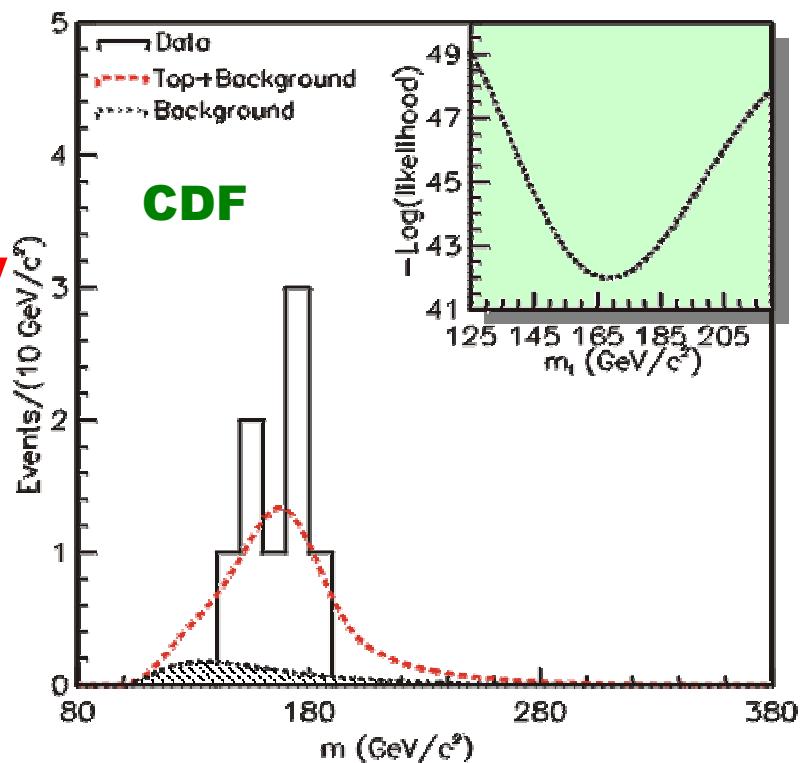
PRD 60, 05001 (1999)



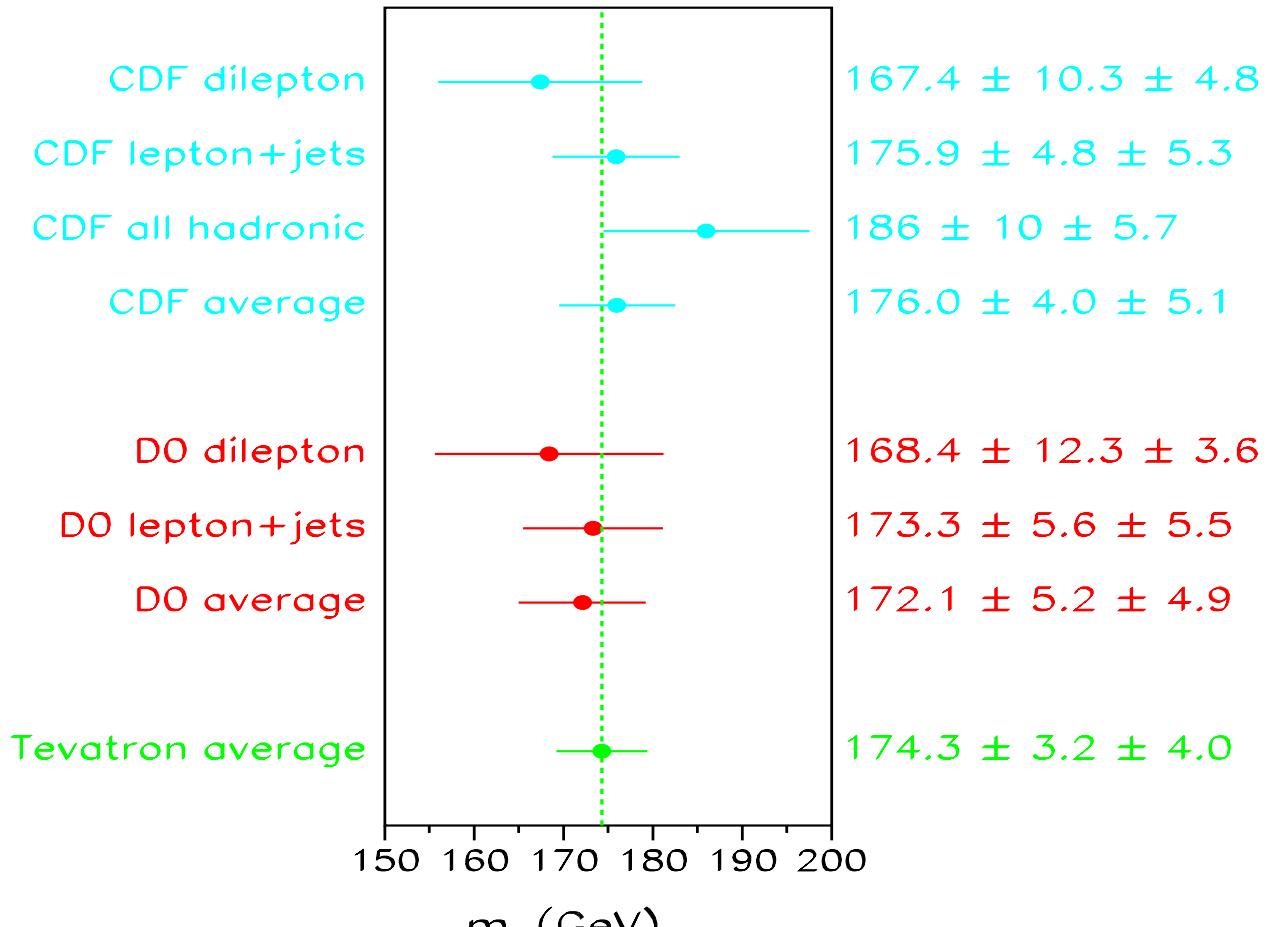
167.4±10.3±4.8 GeV

Largest systematics

Jet energy	3.8 GeV
Gluon radiation	2.7 GeV



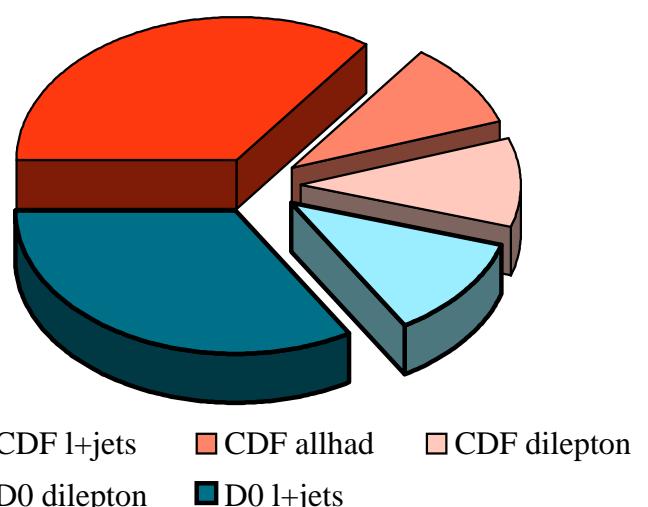
Top Quark Mass Summary



$$m_t = 174.3 \pm 3.2(\text{stat}) \pm 4.0(\text{syst}) \text{ GeV}$$

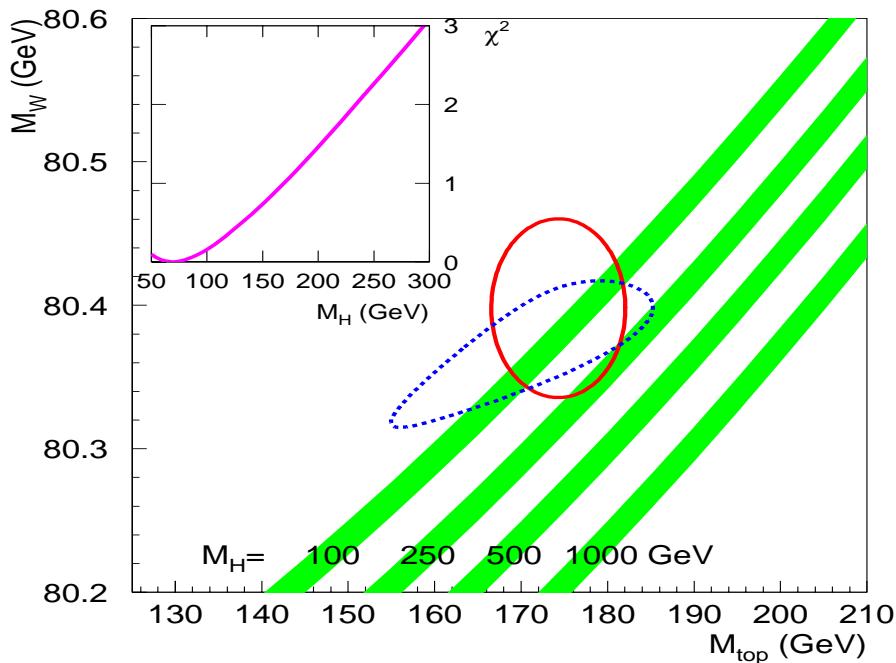
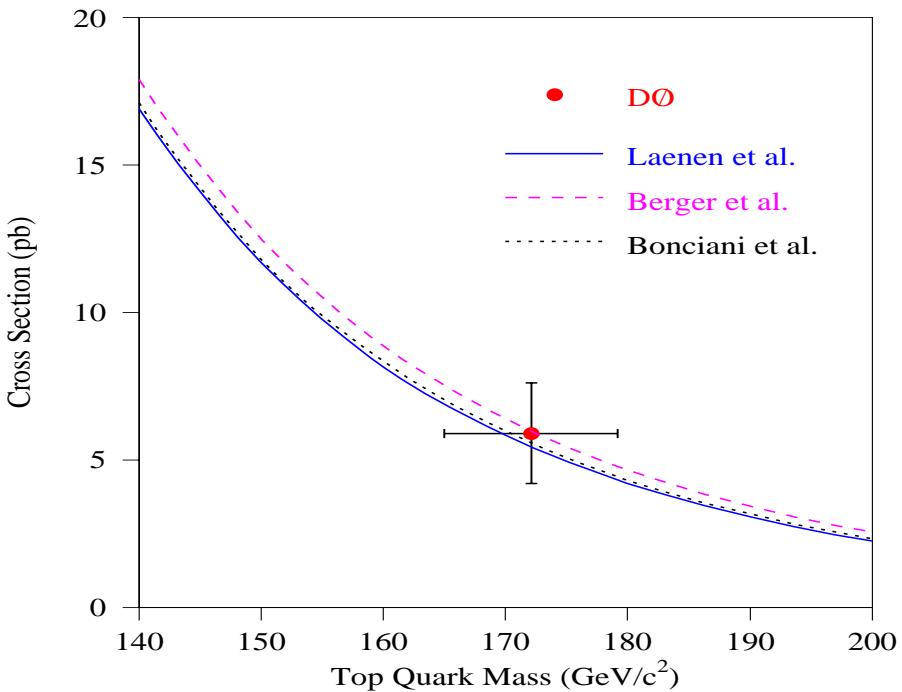
$$= 174.3 \pm 5.1 \text{ GeV}$$

Relative weight in top mass average



Standard Model Comparison

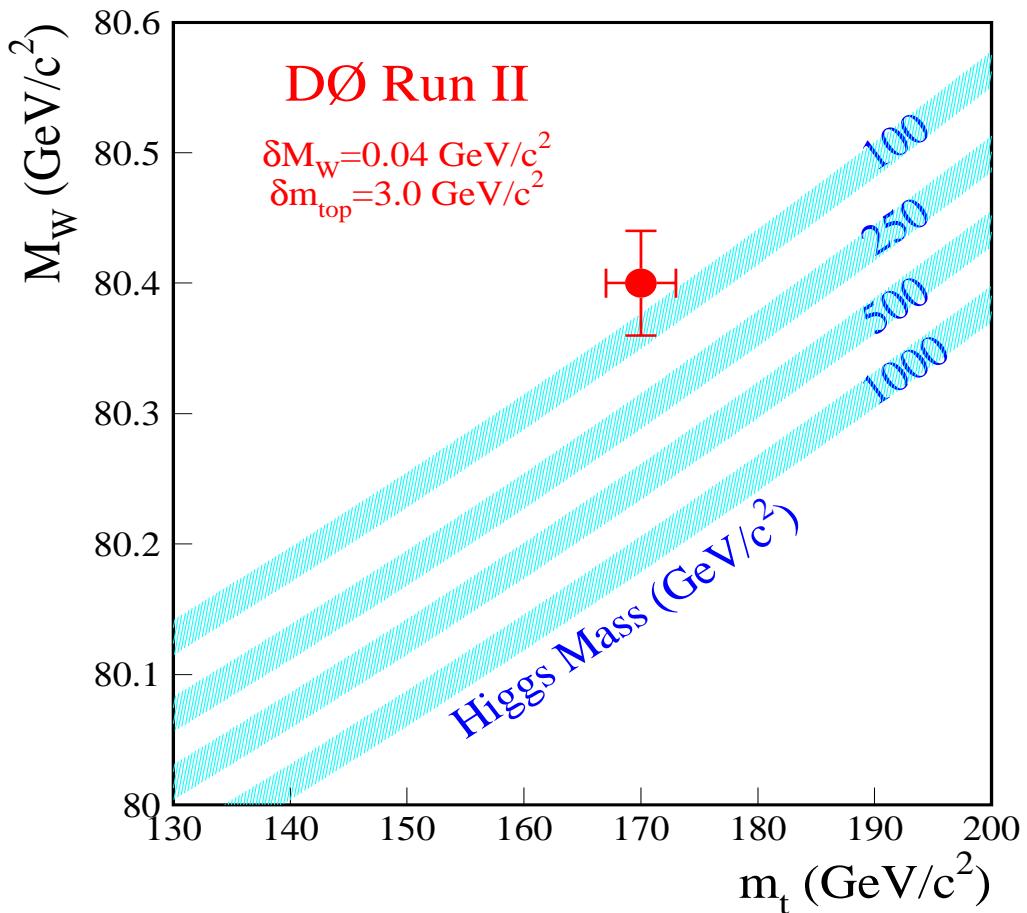
- Good agreement between observation and theoretical calculations
- Data prefer a low Higgs mass



Run II Mass Measurement

Statistical and systematic errors contribute equally to the total errors of the present measurements

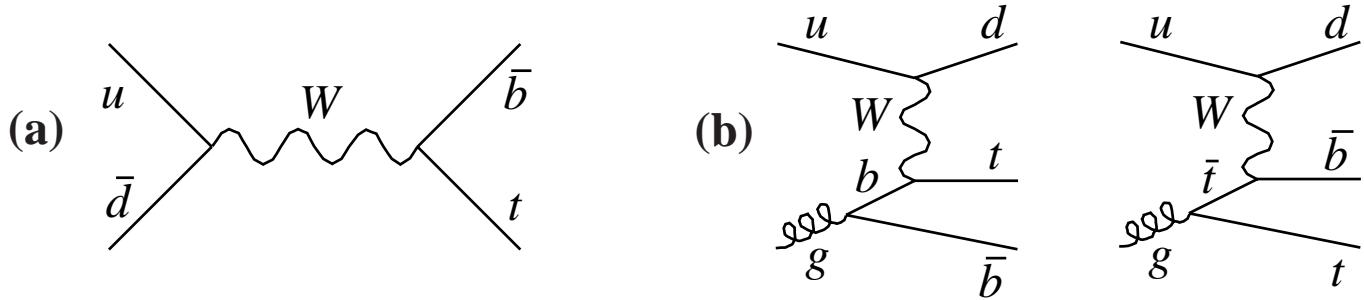
Most of the errors are expected to scale with $1/\sqrt{N}$, the δm_t is expected to reduce to 3 GeV or less



Combined with the data from LEP/SLC, the Higgs mass can be constrained to be within 30%

Single Top Production

Single top quark is produced through electroweak processes



In the standard model

$$\sigma(p\bar{p} \rightarrow tb + X) \propto \Gamma(t \rightarrow Wb) \propto |V_{tb}|^2$$

$$\sigma(W^* \rightarrow tb) = 0.72 \pm 0.10 \text{ pb}$$

$$\sigma(Wg \rightarrow tbq) = 1.7 \pm 0.2 \text{ pb}$$

Measure single top quark cross section

- Direct access to the Wtb vertex
- Directly measure the top quark decay width and CKM matrix element
- Probe anomalous couplings

Final State $W^* \rightarrow tb \rightarrow Wb\bar{b} \rightarrow (\ell\nu)b\bar{b}$
 $Wg \rightarrow tbq \rightarrow Wb\bar{b}q \rightarrow (\ell\nu)b\bar{b}q$

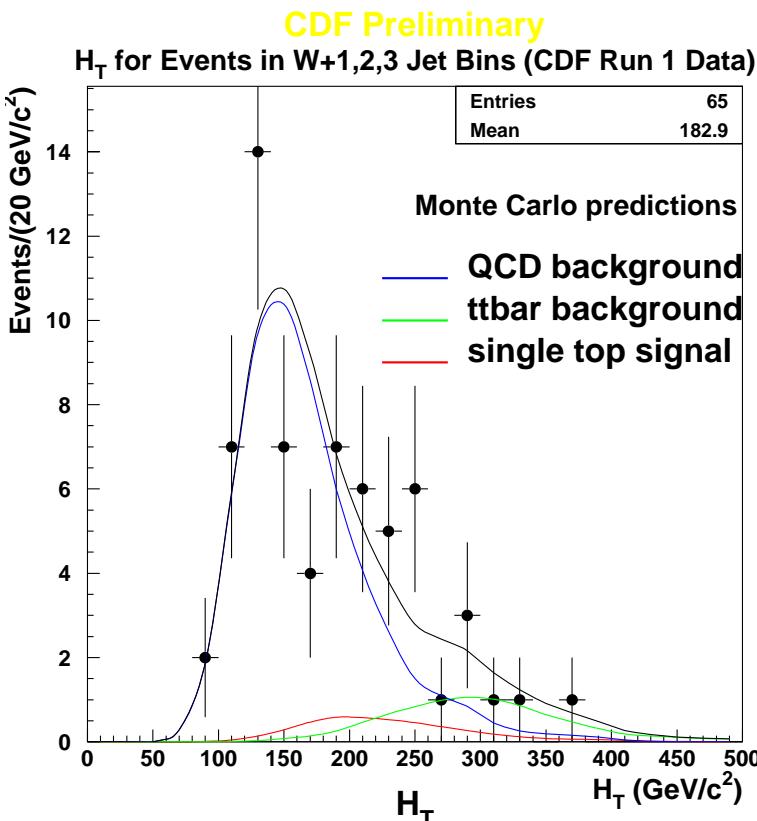
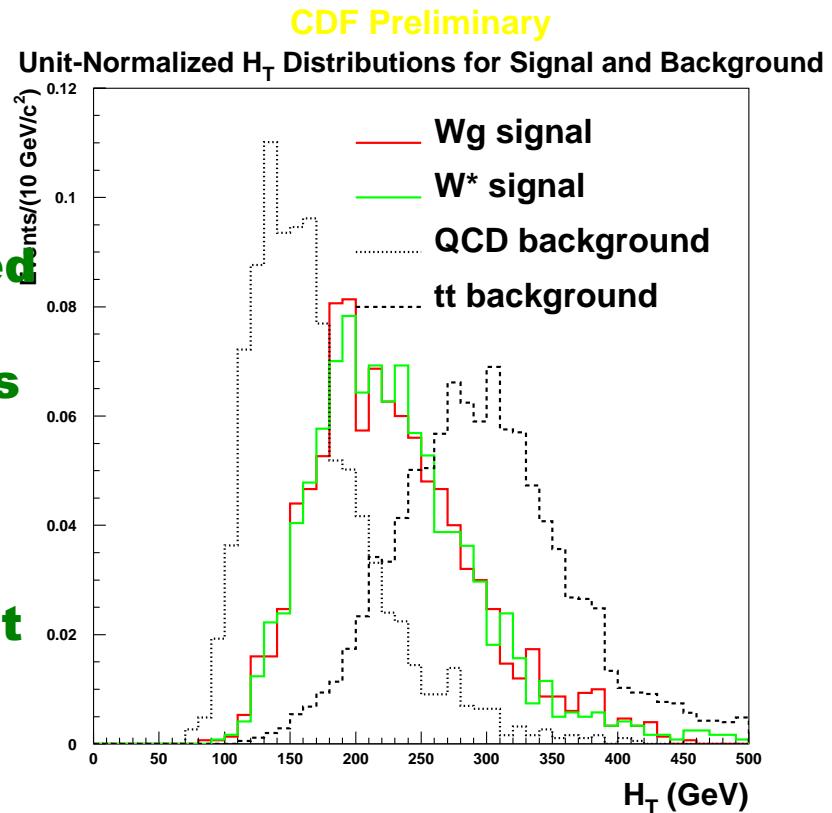
(In Wg mode, one of the b-jets is very soft)

Major backgrounds: $t\bar{t}$ and $Wb\bar{b}$ production

Single Top Production

In many kinematic distributions, the signal is sandwiched between two major background sources

Neither experiment were able to extract a signal in Run I



CDF:

$\sigma < 13.5 \text{ pb} @ 95\% \text{ CL}$

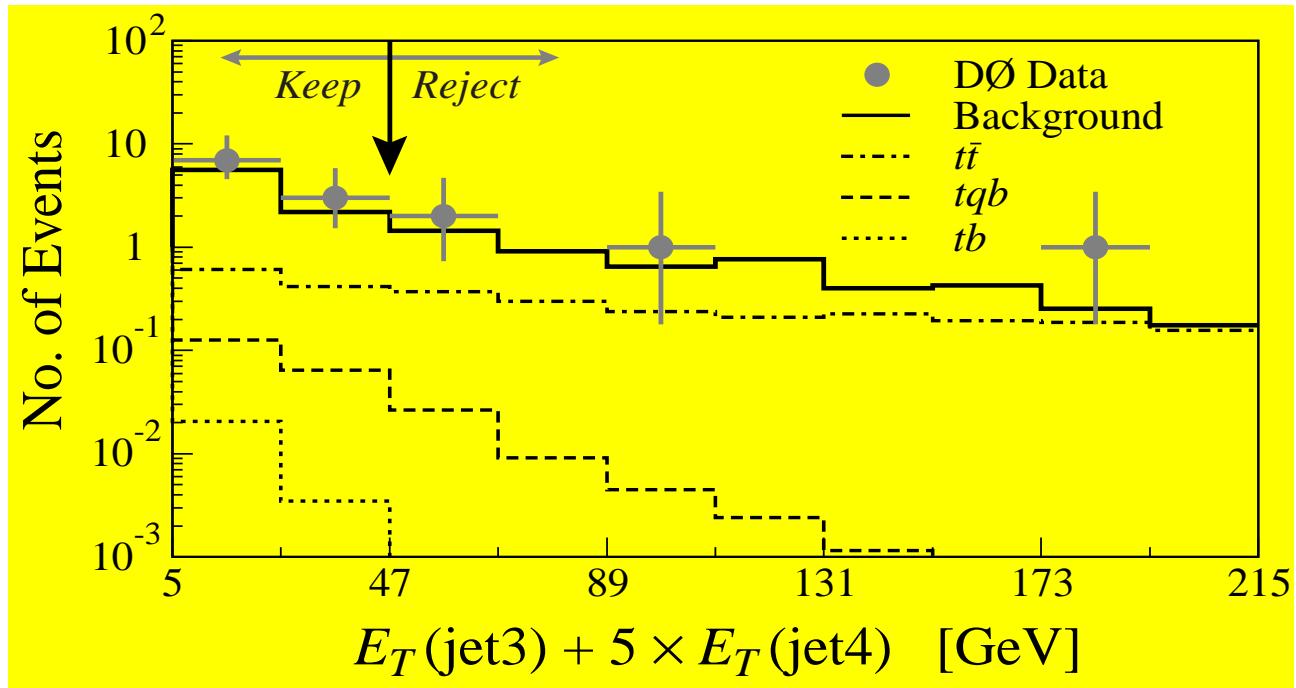
DØ:

$W^* \rightarrow tb: \sigma < 39 \text{ pb}$
 $Wg \rightarrow tbq: \sigma < 58 \text{ pb}$

Single Top Production

Single top search suffers from

- **small production cross section**
- **in Wg production, one of the two b-jets is soft and mostly in the forward region**
- **jet multiplicity is low**
- **inefficient b-jet tagging**



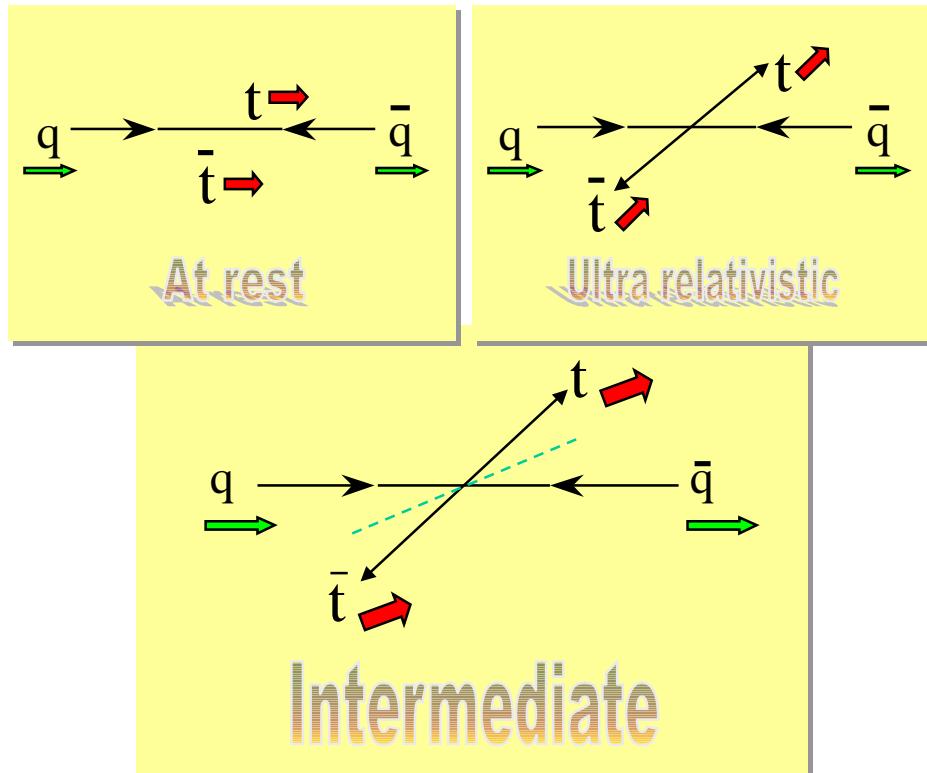
In Run IIa

- **a factor of 20 increase in luminosity**
- **expect to have a factor of 2.5 more from the improved detector and tools**

We expect to observe this process in Run II and measure $\delta\Gamma_t \sim 20\% \Gamma_t$

Spin Correlation

The spins of t and \bar{t} produced from $q\bar{q}$ annihilation are expected to have strong correlations

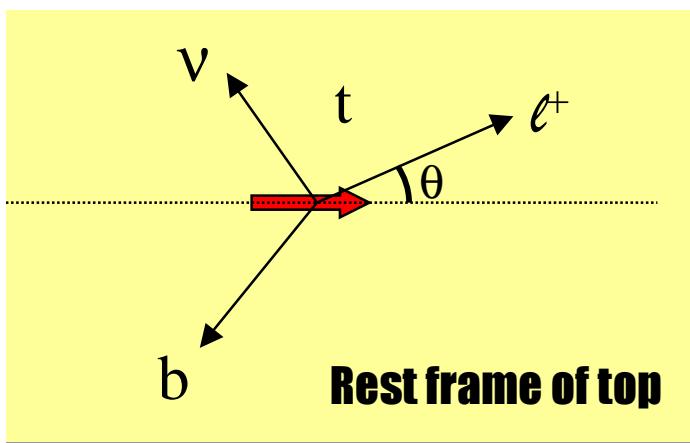


- 1) Since $\Gamma_t = \Gamma(t \rightarrow Wb) \approx 1.5 \text{ GeV} \gg \Lambda_{QCD} \approx 150 \text{ MeV}$ top quarks decay before hadronization
- 2) the polarization of the top quark is carried by its decay products
- 3) At $\sqrt{s} = 1.8 \text{ TeV}$, 90% of $t\bar{t}$ is produced through $q\bar{q}$ annihilation

The top quark polarization can be reconstructed, not possible for other quarks!

Spin Correlation

Charged leptons and down-type quarks are most sensitive to the top quark polarization



$$\frac{1}{\Gamma} \frac{d\Gamma}{d(\cos \theta_i)} = \frac{1 + \alpha_i \cos \theta_i}{2}$$

Particle _i	α_i
l^+ or d	1
ν or u	-0.31
W^+	0.41
b	-0.41

- Down-type quarks are impractical to identify, only dilepton events are considered
- In an optimized spin quantization basis, only like-spin combinations are produced
(G. Mahlon and S. Parke, PLB 411, 173 (1997))
- In this basis, the spin correlation can be expressed using angles

$$\frac{1}{\sigma} \frac{d^2\sigma}{d(\cos \theta_+) d(\cos \theta_-)} = \frac{1 + \kappa \cos \theta_+ \cos \theta_-}{4}$$

correlation parameter
SM value ≈ 0.9

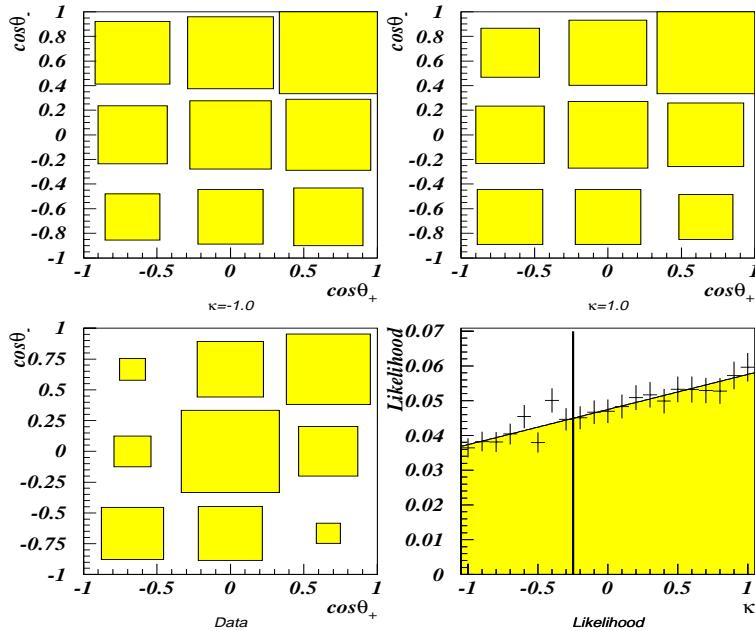
Spin Correlation

- direct probes the properties of ‘bare’ quarks, free of hadronization effect
- probes for non-standard interactions at the production as well as at the decay vertices

Six dilepton events are analyzed assuming $t\bar{t}$ decay hypothesis and $m_t = 175 \text{ GeV}$

Both lepton and jet pairings and multiple Neutrino solutions are considered and weighted

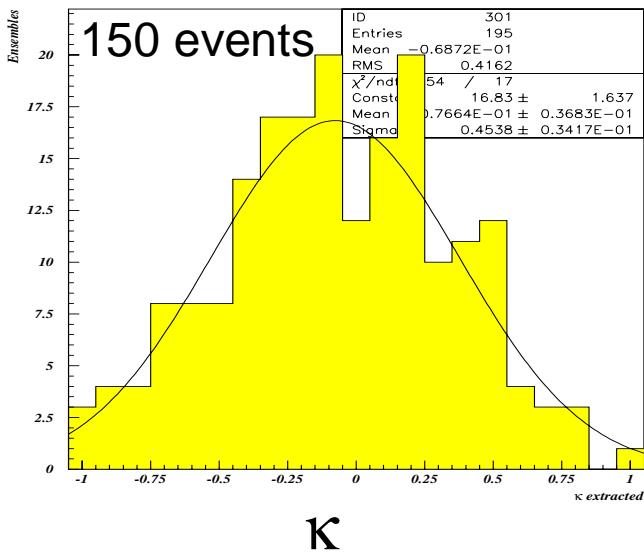
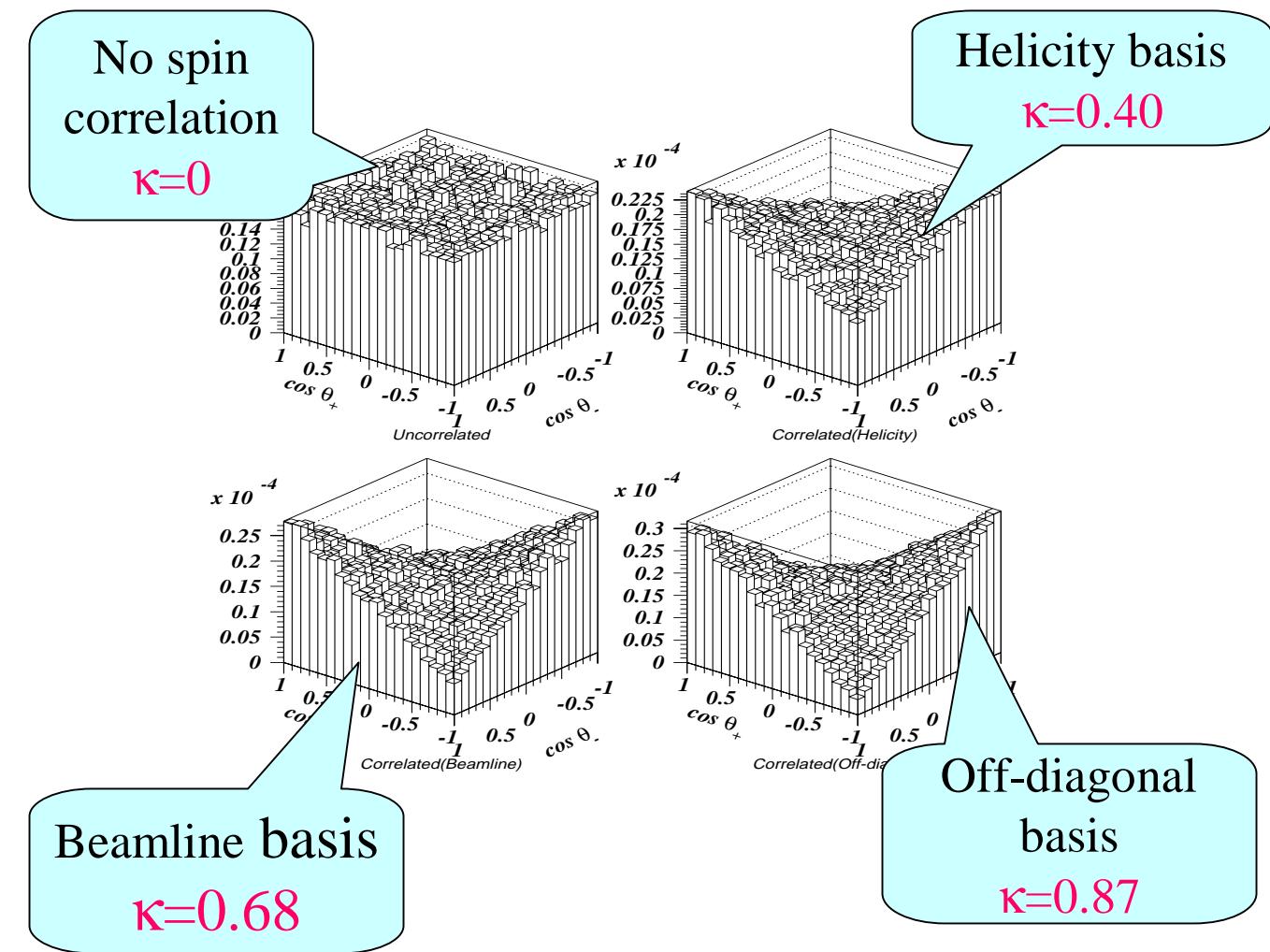
The sensitivity is limited by statistics!



$\kappa > -0.25 @ 68\% \text{ CL}$

DØ: PRL 85, 256 (2000)

Spin Correlation



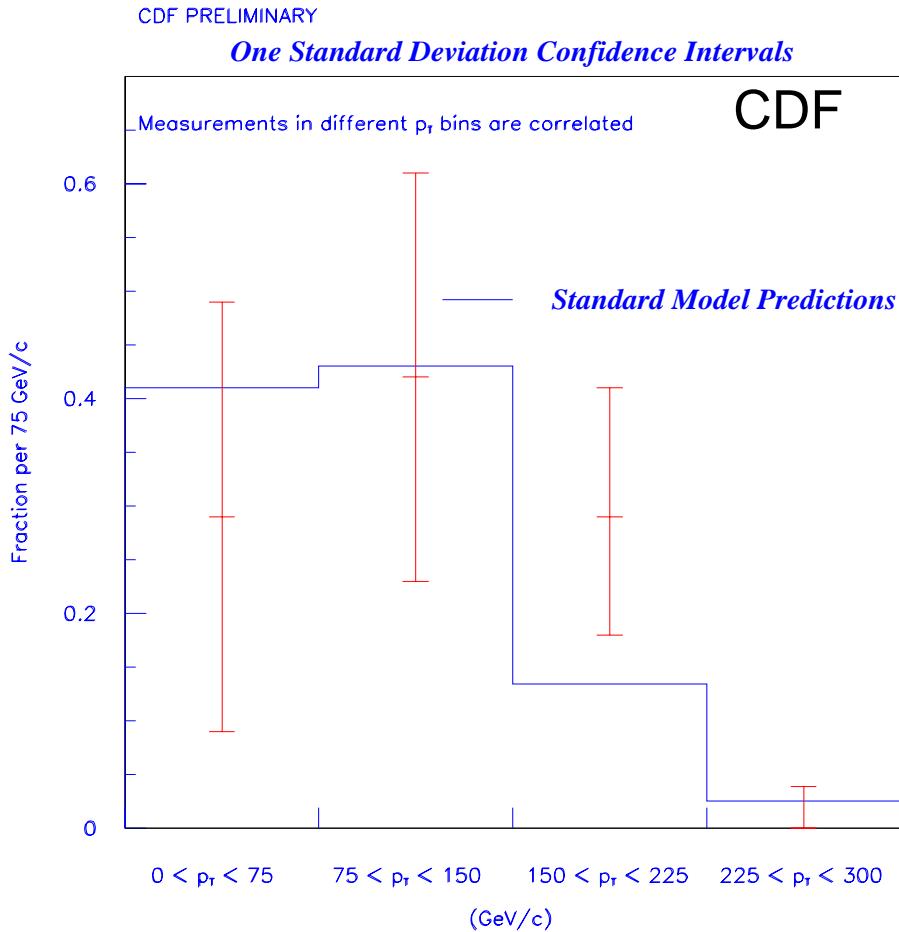
More than 150 dilepton events expected are expected from Run IIa

$\kappa=0$ and $\kappa=1$ can be separated at more than 2σ level

Transverse Momentum

Top Quark Differential Cross Section

$$R_1 + R_2 = 0.72^{+0.13}_{-0.13} (\text{stat})^{+0.06}_{-0.06} (\text{syst})$$
$$R_4 < 0.114 \text{ at } 95\% \text{ C.L.}$$

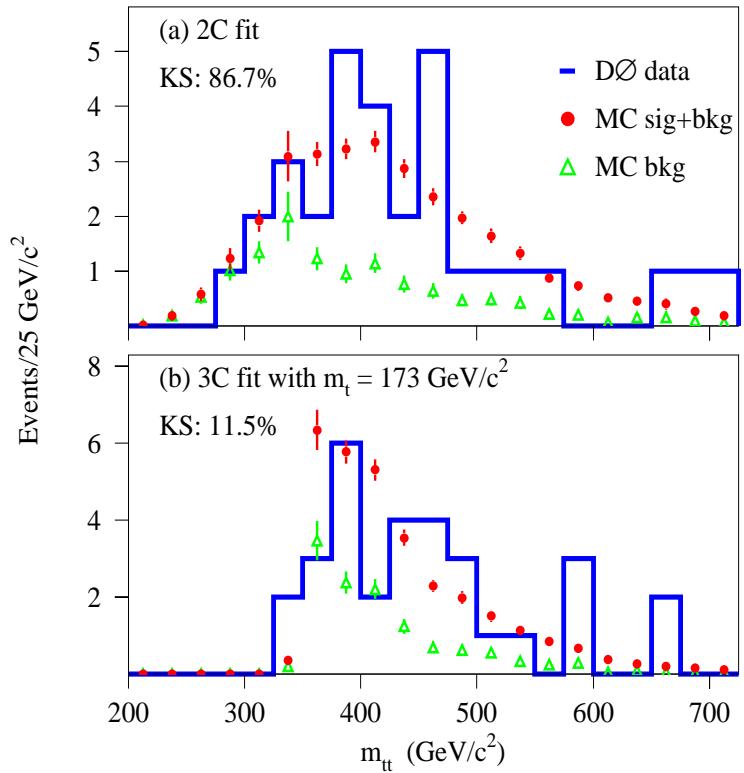
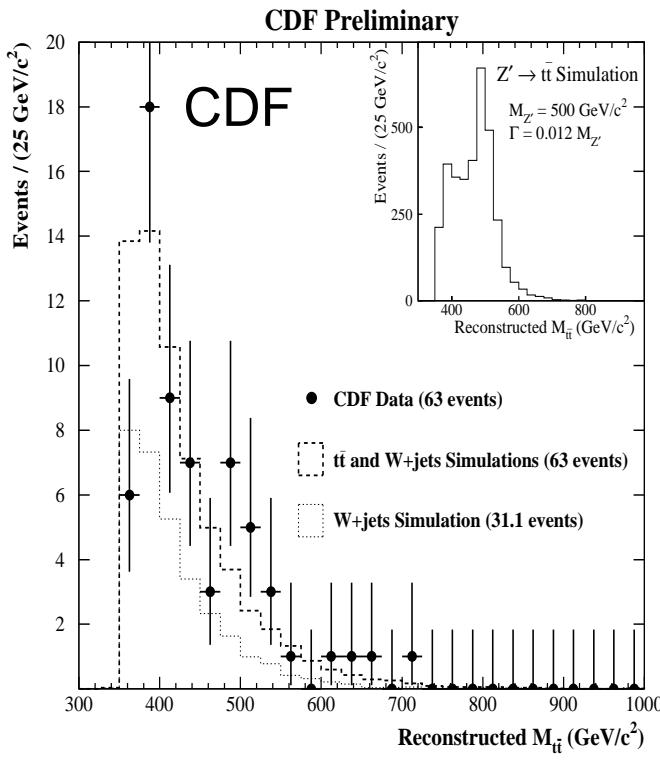


p _T Bin	Measured Fraction of Top Quarks
0 < p _T < 75 GeV/c	$R_1 = 0.29^{+0.18}_{-0.18} (\text{stat})^{+0.08}_{-0.08} (\text{syst})$
75 < p _T < 150 GeV/c	$R_2 = 0.42^{+0.18}_{-0.18} (\text{stat})^{+0.05}_{-0.07} (\text{syst})$
150 < p _T < 225 GeV/c	$R_3 = 0.29^{+0.12}_{-0.10} (\text{stat})^{+0.06}_{-0.05} (\text{syst})$
225 < p _T < 300 GeV/c	$R_4 = 0.000^{+0.035}_{-0.000} (\text{stat})^{+0.019}_{-0.000} (\text{syst})$

Another tool for investigating anomalous tt production mechanisms

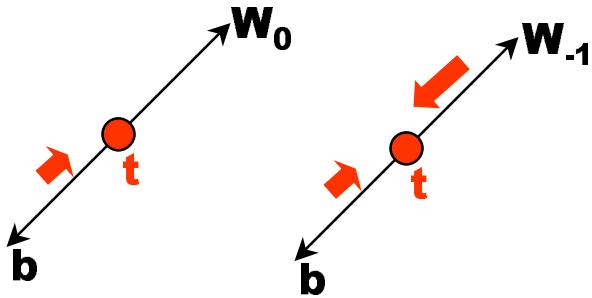
Top-Antitop Mass

- Unlike the bottom quark, top-antitop are not expected to form bound states due to its short lifetime
- However, many models predict the existence of top and antitop resonances
 - technicolor $gg \rightarrow \eta_T \rightarrow (tt, gg)$
 - topcolor $qq \rightarrow V_8 \rightarrow (tt, bb)$
- Top-antitop invariant mass distribution is a tool for searching for new physics



W Boson Helicity

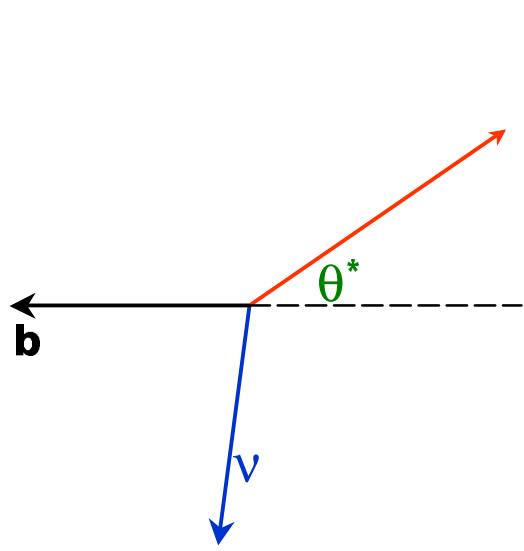
**Top quark decays before hadronization
via V- A interaction**



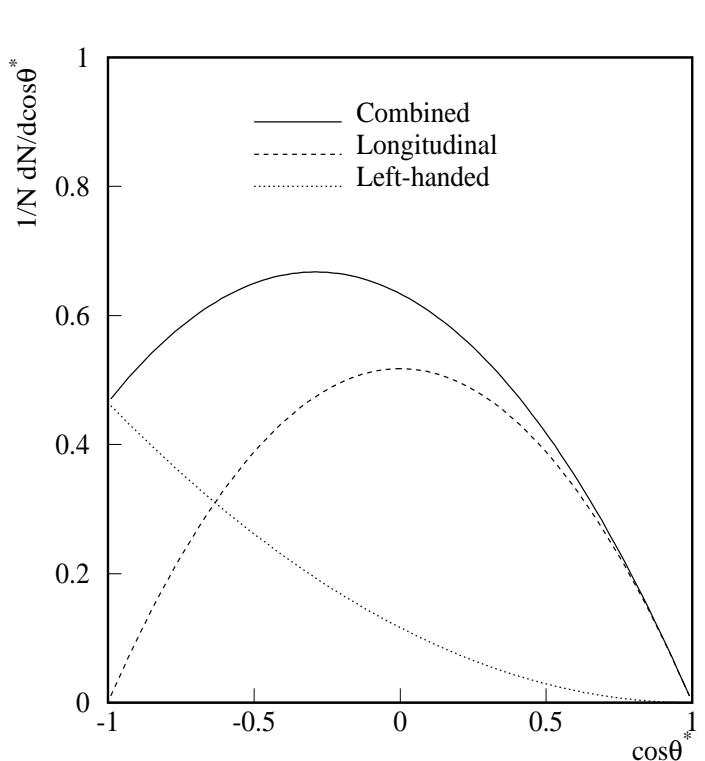
Top quarks can only decay to longitudinal (W_0) or left-hand (W_{-1}) polarized W bosons

In the standard model $F_0 \equiv \frac{W_0}{W_0 + W_{-1}} \approx \frac{m_t^2}{2M_W^2 + m_t^2} \approx 70\%$

W polarizations can be analyzed from the angular or pT distributions of the charged leptons



Top rest frame



W Boson Helicity

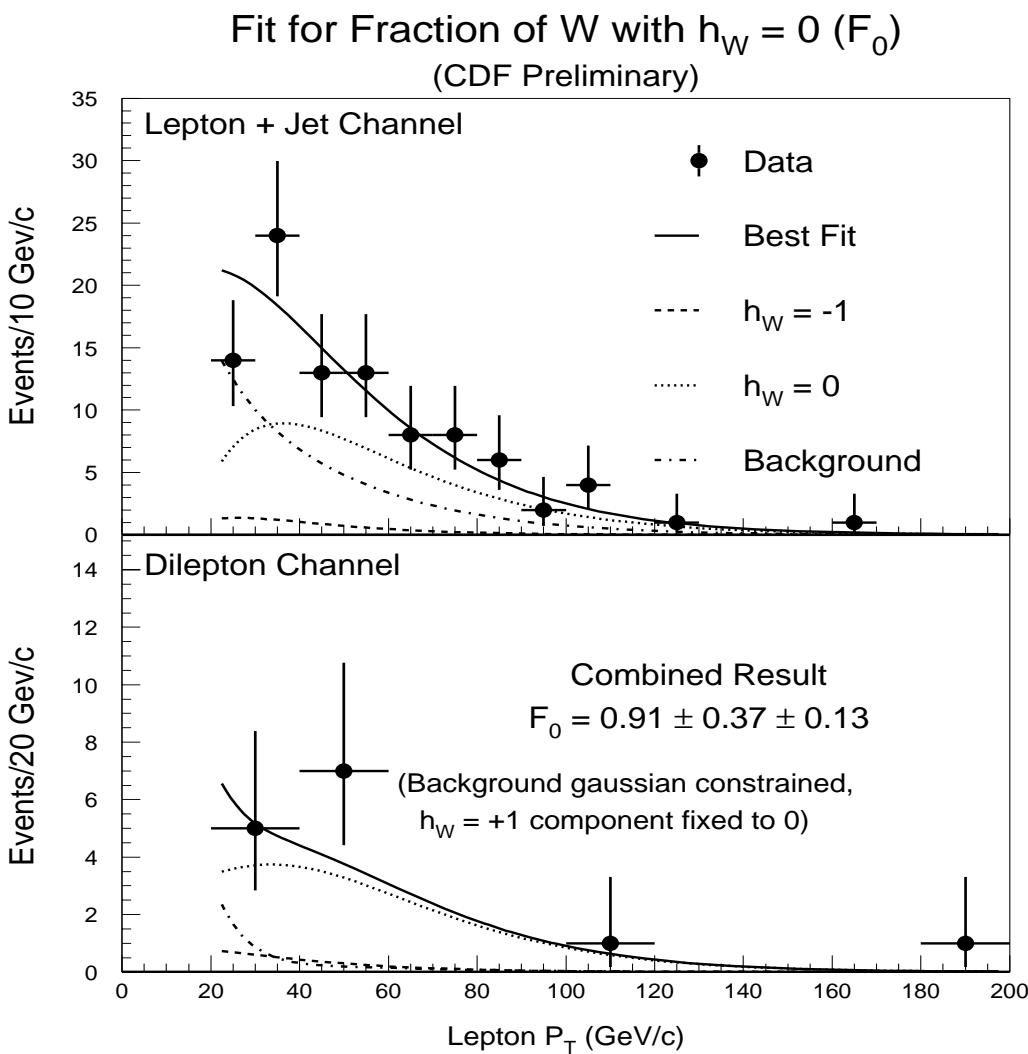
- Top quark decays provide the first opportunity for studying longitudinally polarized Ws
- Non-standard top quark decays may result in different W polarization

**Using dilepton and lepton+jets events,
CDF has measured**

$$F_0 = 0.91 \pm 0.37(\text{stat}) \pm 0.12(\text{syst})$$

$$F_{+1} = 0.11 \pm 0.15(\text{stat}) \pm 0.06(\text{syst})$$

CDF, PRL 84, 216 (2000)



CKM Matrix Element $|V_{tb}|$

- $|V_{td}|^2 + |V_{ts}|^2 + |V_{tb}|^2 = 1$ assuming 3 generations and $|V_{tb}|$ is expected to be very close to 1
- no constraints if more than 3 generations
- any departure of $|V_{tb}|$ from 1 indicates physics beyond the standard model

Extract $|V_{tb}|$ from the measurement of R

$$R = \frac{B(t \rightarrow Wb)}{B(t \rightarrow Wq)} = \frac{|V_{tb}|^2}{|V_{td}|^2 + |V_{ts}|^2 + |V_{tb}|^2}$$

Count dilepton and lepton+jets events with zero, single and double b-tags,
CDF has determined:

$|V_{tb}| > 0.75$ @ 95% CL assuming 3 generations
 $|V_{tb}| > 0.046$ @ 95% CL without the assumption

In Run IIa with 2 fb^{-1} , $\delta|V_{tb}| \approx 10\%$
benefit from large statistics and improved
b-tagging capability

Search for Charged Higgs

- The standard model with two Higgs doublets will result five Higgs particles

$$h^0 \quad H^0 \quad A^0 \quad H^+ \quad H^-$$

- The Higgs sector will have six free parameters
 - $m_{h^0} \ m_{H^0} \ m_{A^0} \ m_{H^\pm}$
 - $\tan \beta = \frac{v_2}{v_1}$: the ratio of VEVs
 - α : Higgs mixing angle

- If charged Higgs bosons are sufficiently light, they can be produced in $t \rightarrow H b$ decays
- The new decay mode will compete with the standard model $t \rightarrow W b$ mode
- Since W and H decay differently

$t \rightarrow H b$ decay will lead to different signatures for top quark pair events

Signature for charged Higgs production:

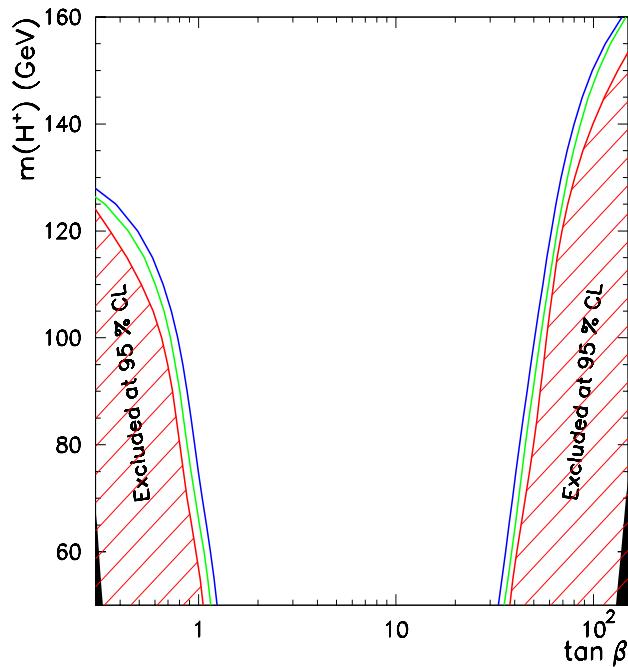
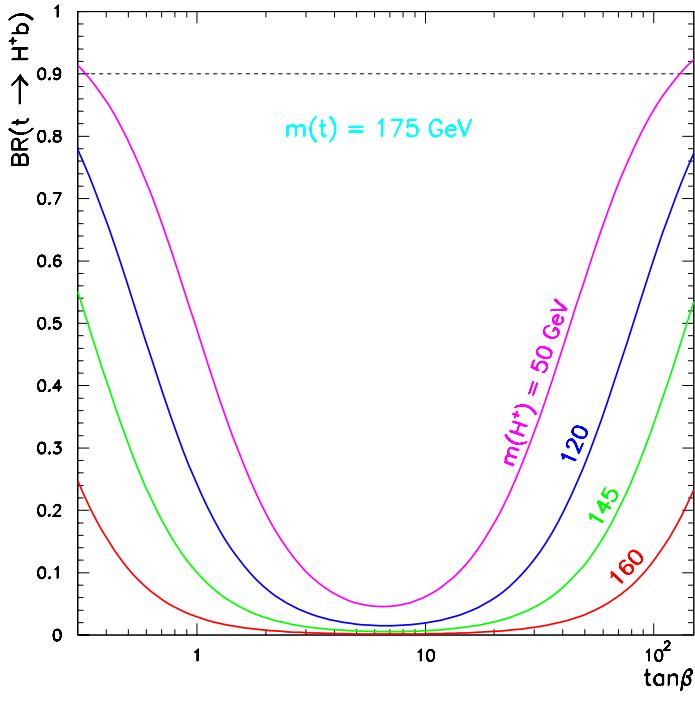
- disappearance of standard $WWbb$ signature
- anomalous τ lepton production

Search for Charged Higgs

Disappearance Search

- The measured top quark pair production cross section agrees with standard model prediction
- Sensitive only to the regions of parameter space with large $B(t \rightarrow H^+ b)$
- Sensitive only to topologies different from WWbb of the SM top quark pair

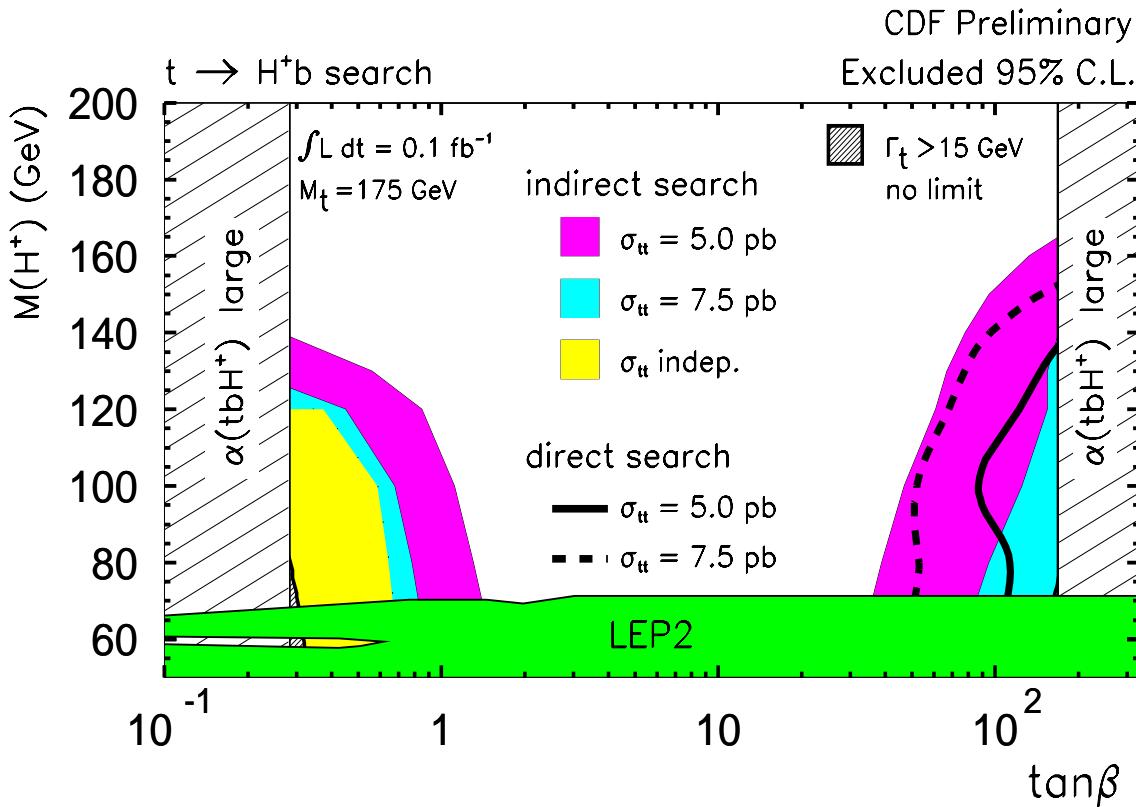
Both CDF and DØ studied the implication of the cross section measurement on the charged Higgs parameter space



Search for Charged Higgs

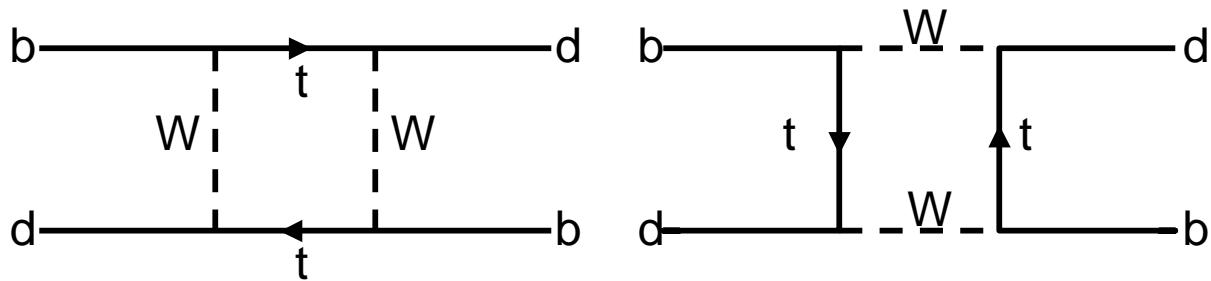
Appearance Search

- CDF searched for $t \rightarrow H^+ b$ decay via τ appearance for high $\tan\beta$ (where $H \rightarrow \tau\nu$)
Phys. Rev. D54, 735 (1996)
- For the τ appearance analyses, τjjX and acoplanar $\tau\tau$ events were searched
- The major backgrounds are fake taus, W+jets, Z+jets and WW, WZ, ZZ productions
- 7 events were observed with 7.4 ± 2.0 events expected. No excess of events



Heavy Top Quark

- Besides negative searches, the other indication of a heavy top quark is from the $B^0\bar{B}^0$ mixing measurement by ARGUS and CLEO



- Theoretically it was known since early 80s that some super - gravity models require top quark mass to be above 150 GeV

Consequences of a heavy top

- Yukawa couplings: $\lambda_t = \frac{\sqrt{2}m_t}{v} \approx 1$ for top quarks and $\lambda_e = \frac{\sqrt{2}m_e}{v} \approx 3 \times 10^{-6}$ for electrons
- Generates electroweak symmetry breaking through radiative correction

Is it “Why the top quark is so heavy ?” Or
Is it “Why other particles are so light ?”

Top Physics in Run II

**Expect to have a factor of 50 increase
in top quark sample in Run IIa**

- Improved measurement of top quark mass and pair production cross section
- Measure top quark decay branching ratios and test of standard model predictions
- Study kinematic distributions
 - spin correlation
 - top polarization
 - transverse momentum
 - mass of top quark pair
 -
- Observation of single top quark production and determination of its decay width
- Search for new physics from top quark decays

**Hope to answer the question
“Does the top quark play a role
in electroweak symmetry breaking?”**

A rich top physics program!